Phenomenology of Hard Diffraction

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*** in collaboration with M. M. Machado and G. G. Silveira**

Outline

- Electroweak Vector boson processes
  - W^{+-} and Z^0 production
- Quarkonium hadroproduction at NLO
  - Application to Heavy-Ion Collisions
- Quarkonium production in NRQCD factorization
  - J/psi + gamma
  - Upsilon + gamma
  - Nuclear production
- Higgs boson production
  - Diffractive factorization
  - Ultraperipheral Collisions

Regge Theory and Pomeron

✓ Ressonances as observables in *t* channel → meson exchange

✓ t channel trajectory → ressonances with same quantum numbers (Reggeons)

$$\alpha(t) = \alpha(0) + \alpha' t$$

slope

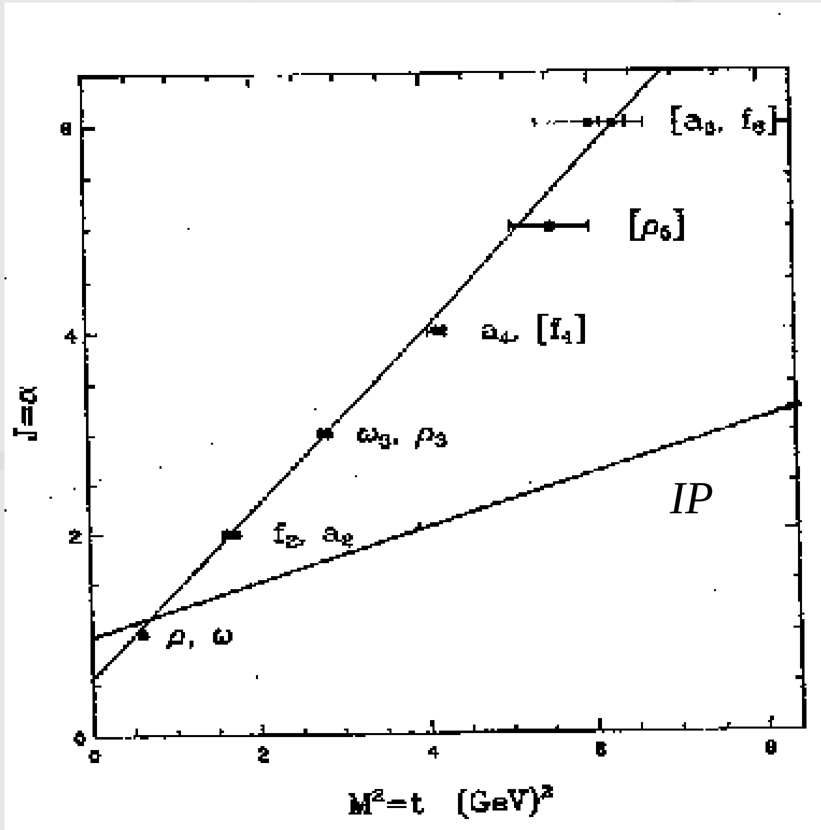
INCREASE AT HIGH ENERGIES

❖ Chew and Frautschi (1961) and Gribov (1961) introduced a Regge trajectory with intercept 1 for asymptotic total cross sections

❖ This reggeon was named Pomeron (*IP*)

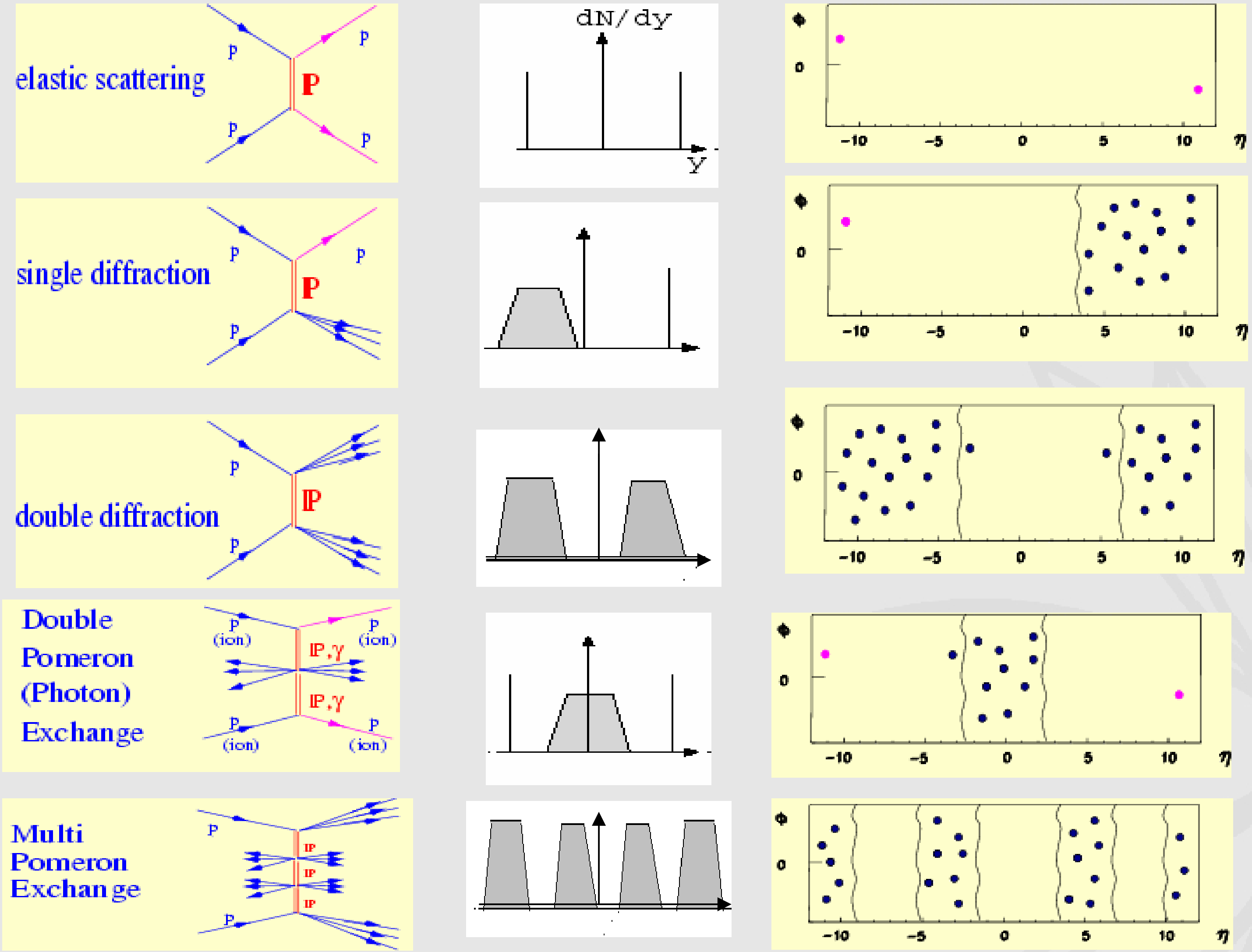
Soft Pomeron values
$\alpha\ (0) \sim 1.09$
$\alpha' \sim 0.25$

P = +1
C = +1
I = 0



α(0) { energy dependence of the diffractive cross section

Diffractive processes

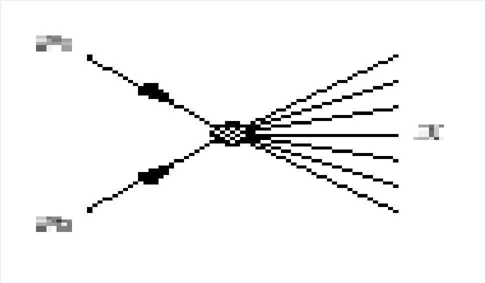


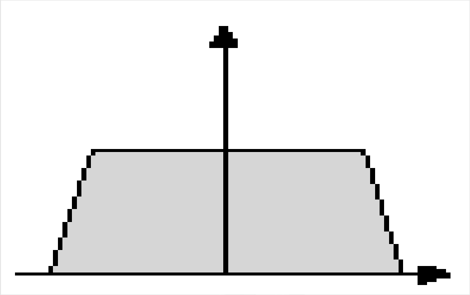
Tevatron/LHC
Higgs: NLO
W, Z
QQ: NRQCD, NLO

Tevatron/LHC
Higgs: photo-,NLO
QQ: NLO

Rapidity

Inelastic scattering





Rapidity

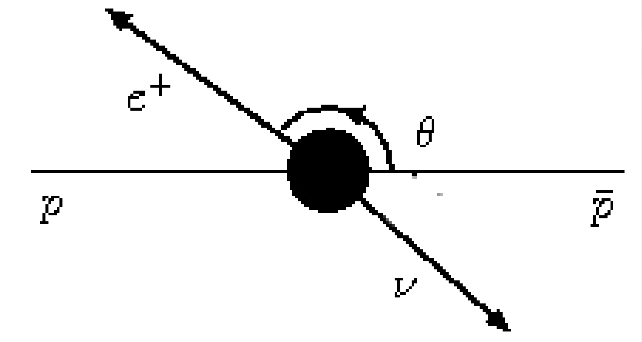
$$y = \frac{1}{2}\ln\frac{E+p_z}{E-p_z} \approx -\ln\tan\frac{\theta}{2} = \eta$$

η ➡ pseudorapidity for a particle with $(E, \vec{p}_\perp, p_z)$ and polar angle θ

Diffraction defined by

➡ leading proton

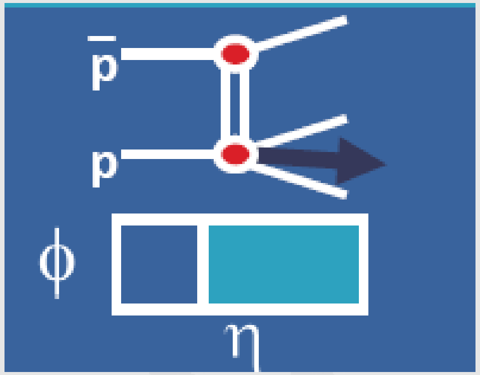
➡ large rapidity gap



Hard Single diffraction

- Large rapidity gap

- Intact hadrons detected



- Diffractive production of some objects is possible to be studied

Jets, W, J/ψ, b ...

- Measurement of the ratio of diffractive to non-diffractive production

Hard component	Fraction (R) %
Dijet	0.75 ± 0.10
W	1.15 ± 0.55
b	0.62 ± 0.25
J/ψ	1.45 v 0.25

All fractions

~ 1%

Goulianos Low x 2009

Diffractive dijet cross section

$$\sigma(\bar{p}p \rightarrow \bar{p}X) \approx F_{jj} \otimes F_{jj}^{D} \otimes \hat{\sigma}(ab \rightarrow jj)$$

❖ Study of the diffractive structure function

$$F_{jj}^{D} = F_{jj}^{D}(x, Q^2, t, \xi)$$

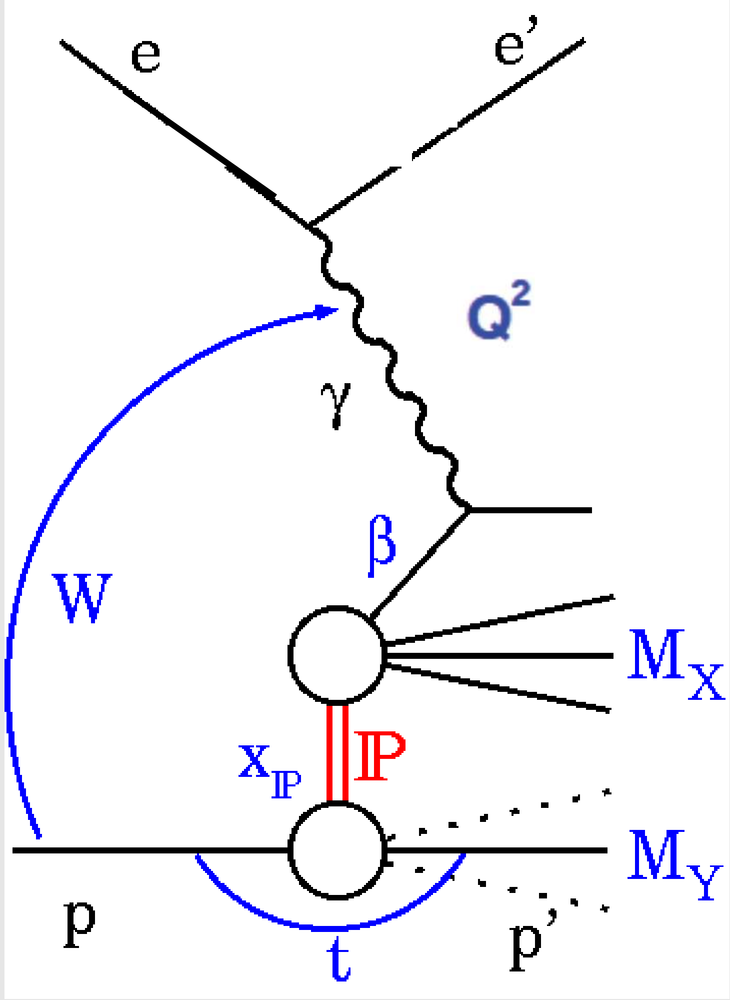
❖ Experimentally determine diffractive structure function

$$R_{\frac{SD}{ND}}(x,\xi) = \frac{\sigma(SD_{jj})}{\sigma(ND_{jj})} = \frac{F_{jj}^{D}(x,Q^2,\xi)}{F_{jj}(x,Q^2)}$$

DATA

KNOWN PDF

Kinematics of DDIS



DDIS

✓ Described by 5 kinematical variables

➢ Bjorken's x

$$x = \frac{Q^2}{2p.q}$$

➢ Squared momentum transfer at the lepton vertex

$$Q^2 = -q^2 = -(k-k')^2$$ or $$y \approx \frac{Q^2}{xs}$$

$$t = -(p'-p)^2$$

$$x_{IP} = \frac{M^2 + Q^2}{W^2 + Q^2}$$

M^2 is the invariant mass of the X system

➢ β is the momentum fraction of the parton inside the Pomeron

$$\beta = \frac{Q^2}{M^2 + Q^2}$$

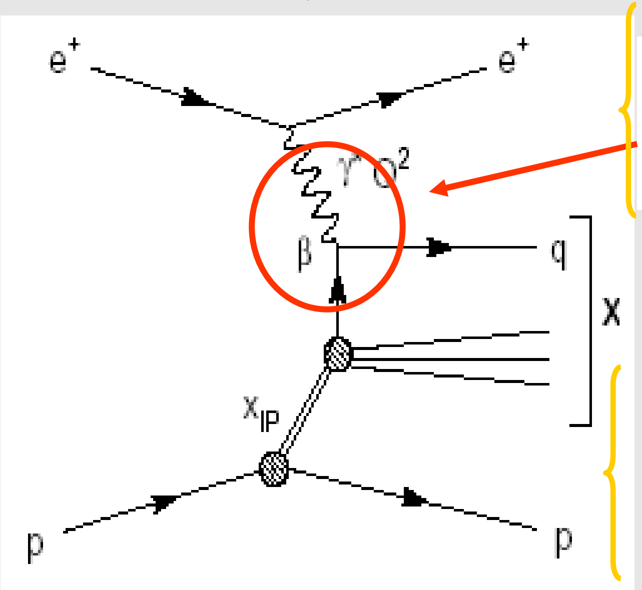
Ingelman-Schlein Model

- IS paper (1985) → first discussion of high-p_T jets produced via Pomeron exchange
- Pomeron → vacuum quantum numbers and with substructure
- Events containing two jets of high transverse energy and leading proton observed in pp scattering at $\sqrt{s} = 630\text{GeV}$ (CERN UA8 experiment, Bonino et al. 1988)
- Rate of jet production in this scattering → 1 – 2%
- Agreement with the predicted order of magnitude made by IS
- Hard diffraction in pp scattering → CDF/D0 Collaborations (Tevatron)
- UA8 group → evidence for a hard Pomeron substructure (Brand et al. 1992)

QCD factorization

PDFs from inclusive diffraction predict cross sections for exclusive diffraction

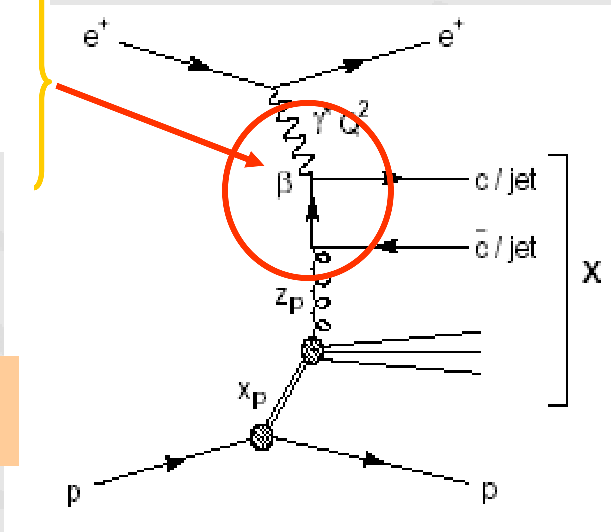
inclusive



hard scattering QCD matrix element perturbatively calculated

Diffractive parton densities are the same for all processes

Exclusive (dijet)



$$\sigma^D(\gamma^* p \to Xp) = \sum_i f_i^D(x, Q^2, x_{IP}, t) \cdot \sigma^{\gamma^* i}(x, Q^2)$$

$\sigma^{\gamma^* i}$ → universal hard scattering cross section (same as in inclusive DIS)

f_i^D → diffractive parton distribution functions → obey DGLAP
universal for diffractive *ep* DIS (inclusive, di-jets, charm)

Diffraction at Tevatron

- Experiments have been investigating diffractive reactions
- First results to diffractive events were reported in 1994-1995
(Abachi et al. 1994; Abe et al. 1995)
- **3** different classes of processes were investigated at the Tevatron

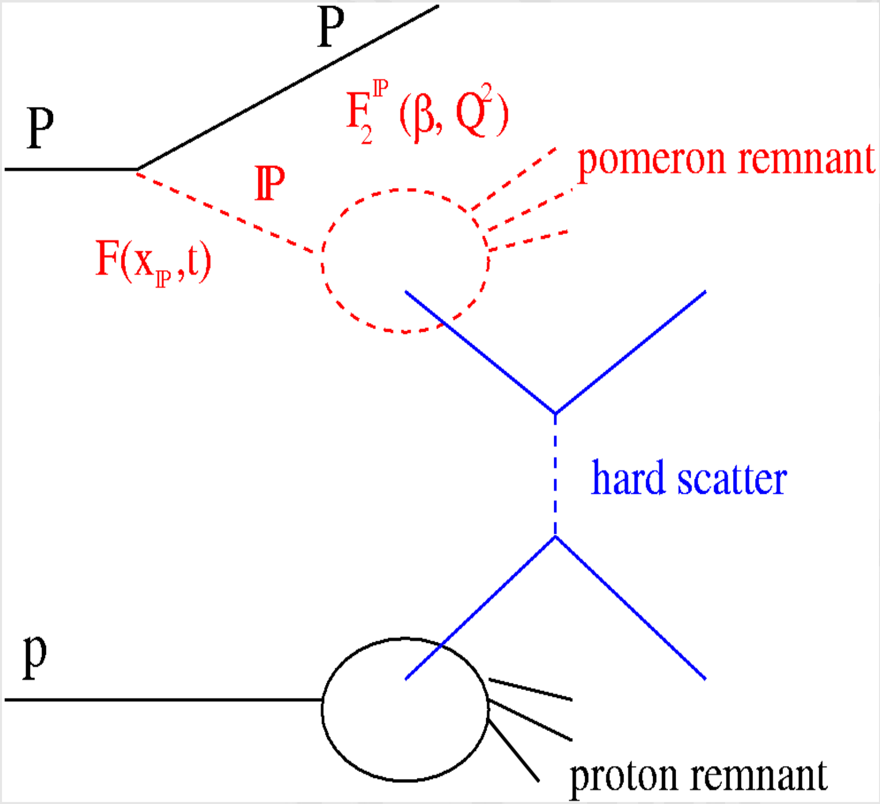
Double diffraction

Single diffraction

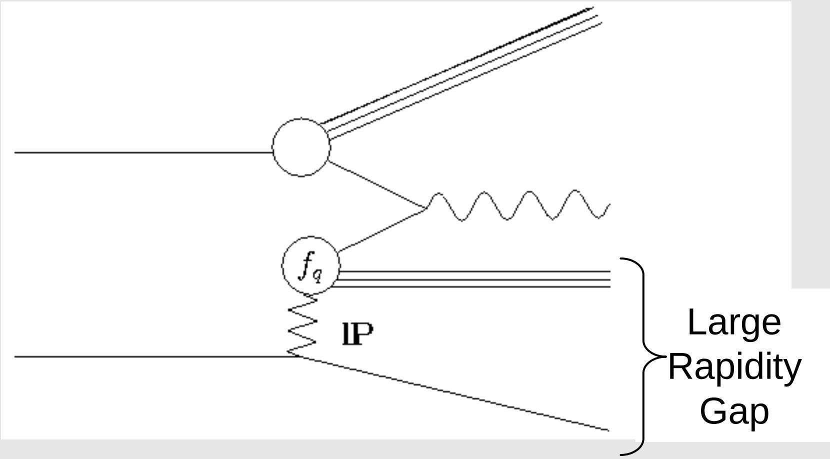
Double Pomeron Exchange

Structure function of the Pomeron

$F_2^{IP}(\beta, Q^2)$



Gap Survival Probability (GSP)



GAP — region of angular phase space devoid of particles

Survival probability — fulfilling of the gap by hadrons produced in interactions of remanescent particles

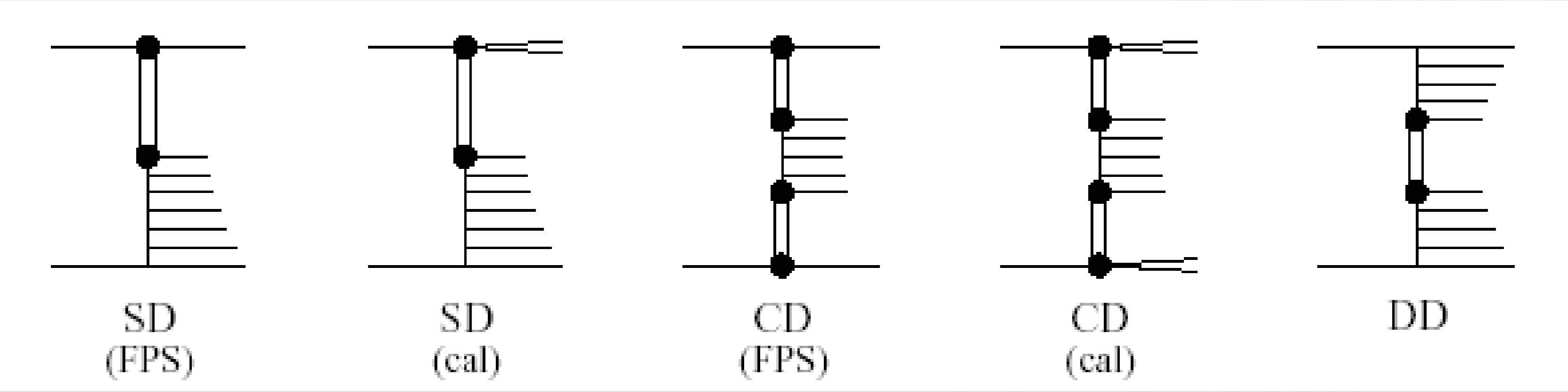
$$<|S|>^2 = \frac{\int d^2b\,|A(s,b)|^2\, P^s(s,b)}{\int d^2b\,|A(s,b)|^2}$$

- *A(s,b)* ➡ amplitude of the particular process (parameter space b) of interest at center-of-mass energy $\sqrt{s}$
- *P^S(s,b)* ➡ probability that no inelastic interaction occurs between scattered hadrons

KMR – Gap Survival Probability

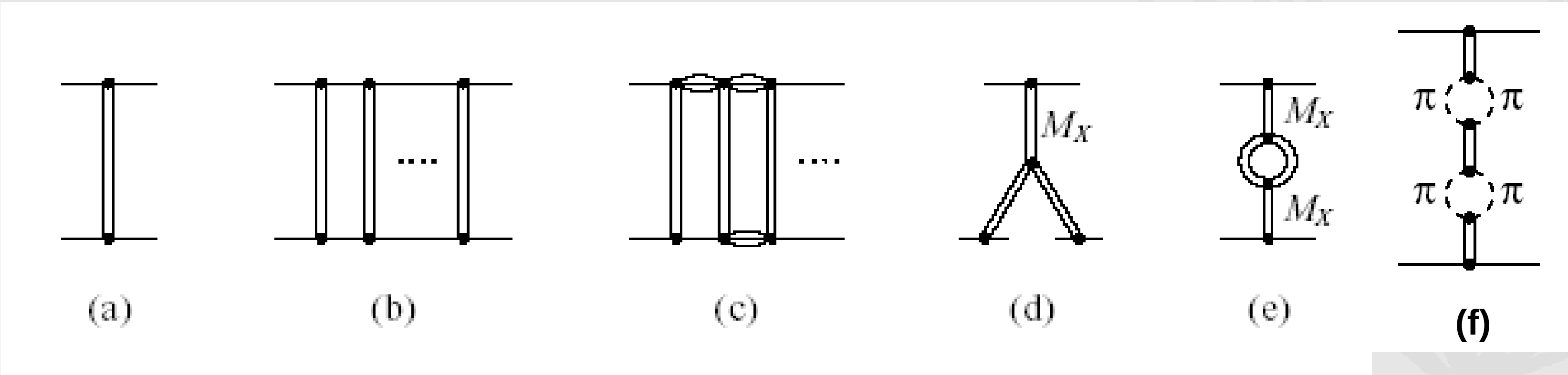
Khoze-Martin-Ryskin Eur. Phys. J. C. 26 229 (2002)

- Survival probability of the rapidity gaps
- Associated with the Pomeron (double vertical line)
  - Calculated
    - * single diffraction (SD)
    - * central diffraction (CD)
    - * double diffraction (DD)
- FPS (cal) → forward photon spectrometer (calorimeter),
- Detection of isolated protons (events where leading baryon is either a proton or a N*)



KMR model

- t dependence of elastic pp differential cross section in the form exp (Bt)
- Pion-loop insertions in the Pomeron trajectory
- Non-exponential form of the proton-Pomeron vertex β (t)
- Absorptive corrections, associated with eikonalization



- (a) Pomeron exchange contribution;
- (b-e) Unitarity corrections to the pp elastic amplitude.
- (f) Two pion-loop insertion in the Pomeron trajectory

KMR model

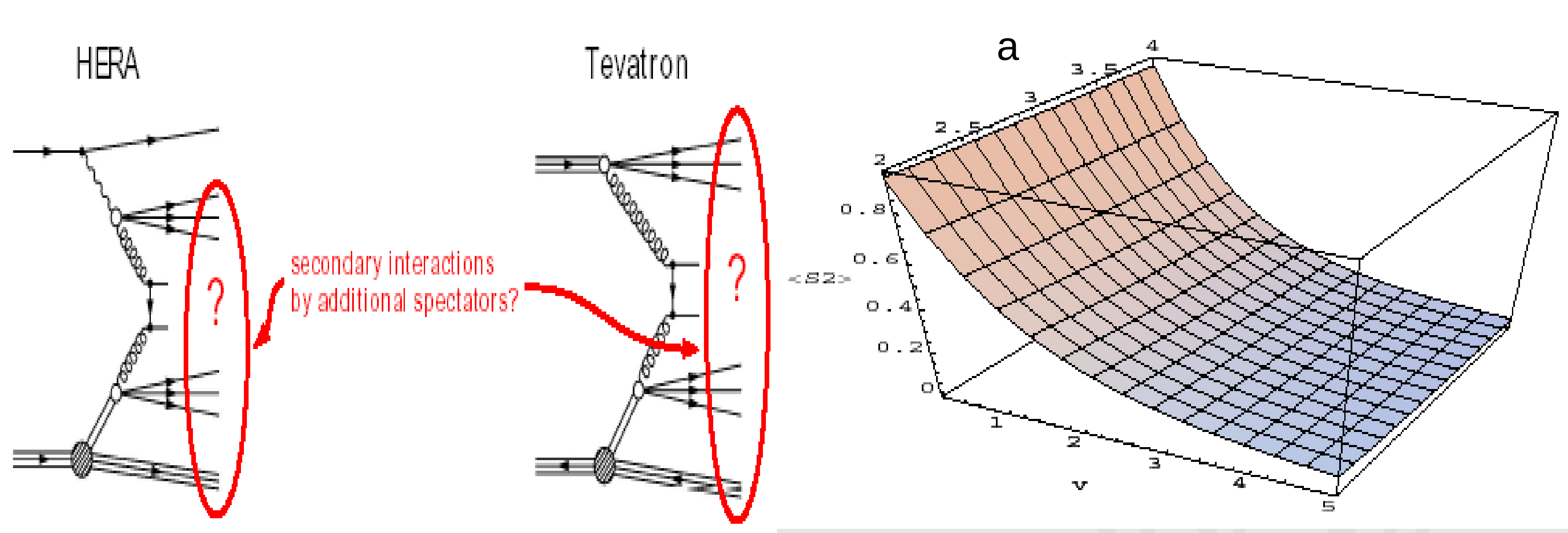
- GSP KMR values

$\sqrt{s}$ (TeV)	$2b$ (GeV^{-2})	Survival probability S^2 for: SD (FPS)	SD (cal)	CD (FPS)	CD (cal)	DD
0.54	4.0	0.14	0.13	0.07	0.06	0.20
	5.5	0.20	0.18	0.11	0.09	0.26
	7.58	0.27	0.25	0.16	0.14	0.34
1.8	4.0	0.10	0.09	0.05	0.04	0.15
	5.5	0.15	0.14	0.08	0.06	0.21
	8.47	0.24	0.23	0.14	0.12	0.32
14	4.0	0.06	0.05	0.02	0.02	0.10
	5.5	0.09	0.09	0.04	0.03	0.15
	10.07	0.21	0.20	0.11	0.09	0.29

- GSP considering multiple channels

GLM - GSP

Gotsman-Levin-Maor PLB 438 229 (1998 - 2002)



- Survival probability as a function of Ω (s,b = 0)
- Ω → opacity (optic density) of interaction of incident hadrons
- Ratio of the radius in soft and hard interactions → a = R_s / R_h
- Suppression due to secondary interactions by additional spectators hadrons

GLM model

GLM - arXiv:hep-ph/0511060v1 6 Nov 2005

- Eikonal model originally → explain the exceptionally mild energy dependence of soft diffractive cross sections

- s-channel unitarization enforced by the eikonal model
- Operates on a diffractive amplitude in different way than elastic amplitude

- Soft input obtained directly from the measured values of σ_{tot}, σ_{el} and hard radius R_H

- *F*1*C* and *D*1*C* → different methods from GLM model

$\sqrt{s}$ (GeV)	$S^2_{CD}(F1C)$	$S^2_{CD}(D1C)$	$S^2_{SD_{incl}}(F1C)$	$S^2_{SD_{incl}}(D1C)$	$S^2_{DD}(F1C)$	$S^2_{DD}(D1C)$
540	14.4%	13.1%	18.5%	17.5%	22.6%	22.0%
1800	10.9%	8.9%	14.5%	12.6%	18.2%	16.6%
14000	6.0%	5.2%	8.6%	8.1%	11.5%	11.2 %

Pomeron flux factor

- x_{IP} dependence is parametrized using a flux factor

$$f_{IP/p}(x_{IP},t) = A_{IP} \frac{e^{B_{IP}t}}{x_{IP}^{2\alpha_{IP}(t)-1}}$$

- IP trajectory is assumed to be linear ⇨ $\alpha_{IP}(t) = \alpha_{IP}(0) + \alpha'_{IP}\, t$

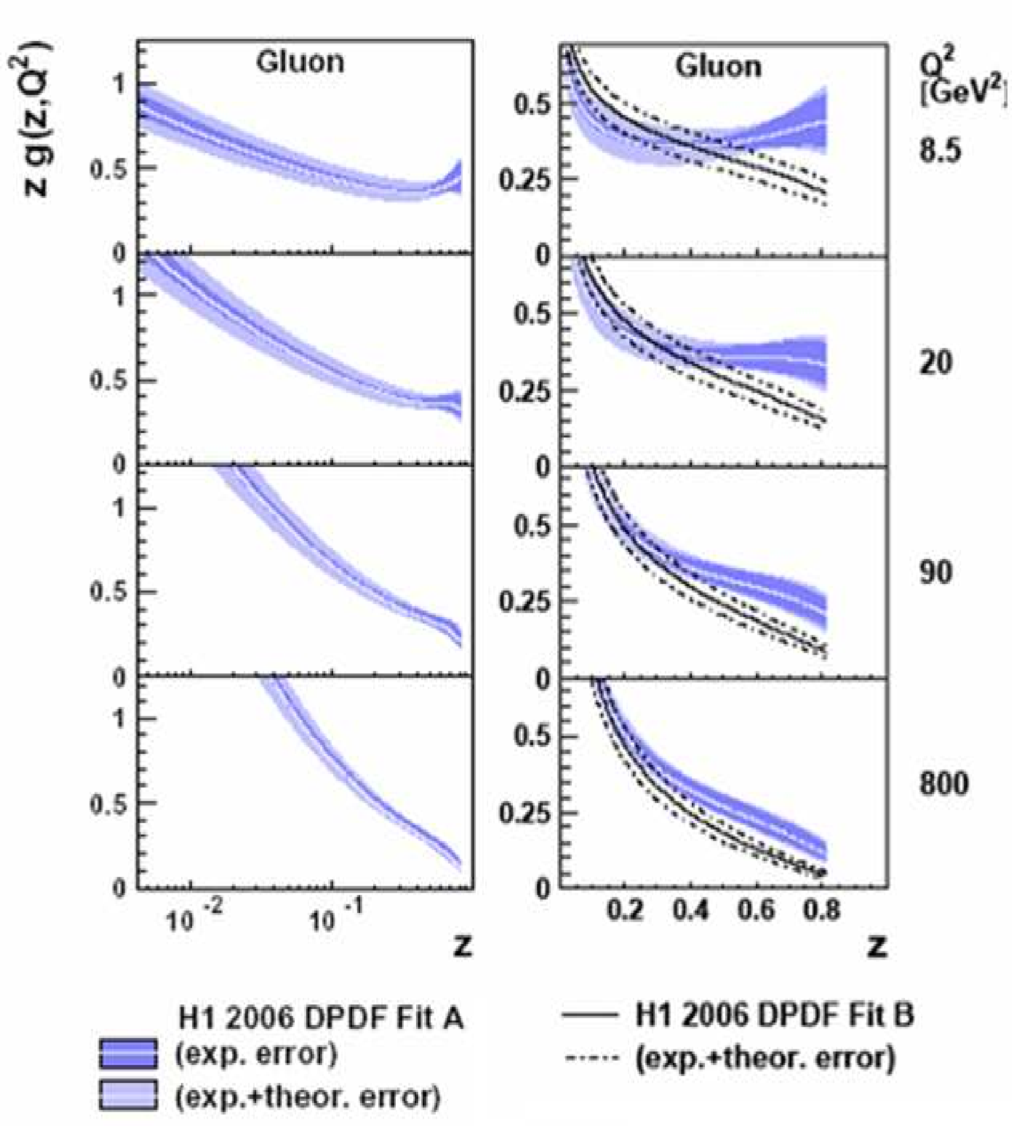
- B_{IP}, α'_{IP} their uncertainties } obtained from the fits to H1 forward proton spectometer (FPS) data

Normalization parameter x_{IP} is chosen such that $x_{IP} \cdot \int_{t_{cut}}^{t_{min}} f_{IP/p}\, dt = 1$ at $x_{IP} = 0.003$

- $|t_{\min}| \approx m_p^2 x_{IP} / (1 - x_{IP})$ is the proton mass
- $|t_{cut}| = 1.0$ GeV^2 is the limit of the measurement

A. Aktas *et al*, Eur. J. Phys. J. C48 (2006) 715

Diffractive Parton Densities (H1-06)



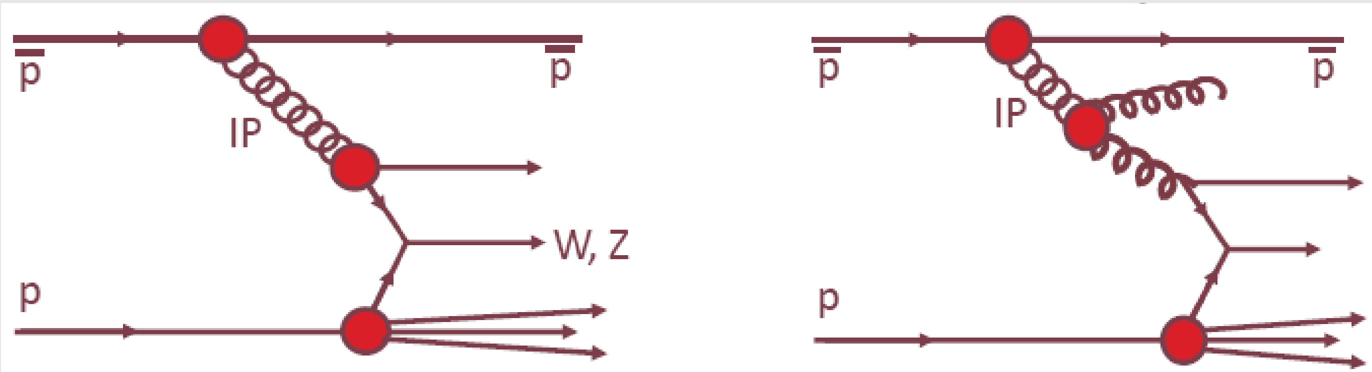
- Total quark singlet and gluon distributions obtained from NLO QCD H1. DPDF Fit A,
- Range 0.0043 < z < 0,8, corresponding to experiment
- Central lines surrounded by inner errors bands → experimental uncertainties
- Outer error bands → experimental and theoretical uncertainties

***z* is the momentum fraction of the parton inner the Pomeron**

Electroweak vector boson production

MBGD, M. M. Machado, M. V. T. Machado, PRD 75, 114013 (2007)

W/Z Production



Leading order ➡ W and Z produced by a quark in the Pomeron

- General cross section for W and Z

$$\frac{d\sigma}{dx_a dx_b} = \sum_{a,b} \int dx_a \, f_{a/p}(x_a, \mu^2) \, f_{b/\bar{p}}(x_b, \mu^2) \frac{d\hat{\sigma}(p\bar{p} \to [W/Z]X)}{d\hat{t}}$$

- W^+ (W^-) inclusive cross section

$$\frac{d\sigma}{d\eta_{e^-(e^+)}} = \sum_{a,b} \int dE_T \, f_{a/p}(x_a) \, f_{b/\bar{p}}(x_b) \left[\frac{V_{ab}^2 G_F^2}{6 s \Gamma_W M_W}\right] \frac{\hat{t}^2 (\hat{u}^2)}{\sqrt{A^2 - 1}}$$

$$\mu^2 = M_W^2 \qquad \hat{t} = -E_T M_W \left[A + \sqrt{(A^2 - 1)}\right]$$

$G_F = 1.166 \times 10^{-5}\ \mathrm{GeV}^{-2}$

- Total decay width ➡ Γ_W = 2.06 GeV
- V_{ab} is the CKM Matrix element
- W^+ (W^-) ➡ dependence in t (u) channel

Energies and Mandelstan Variables

- Total Energy → $E_e = \frac{\sqrt{s}}{4}\left[x_a(1+\cos\theta)+x_b(1-\cos\theta)\right]$
- Longitudinal Energy → $E_L = \frac{\sqrt{s}}{4}\left[x_a(1+\cos\theta)-x_b(1-\cos\theta)\right]$
- Transversal Energy → $E_T = \frac{M_W}{2} sen\theta$
- Mandelstan variables of the process

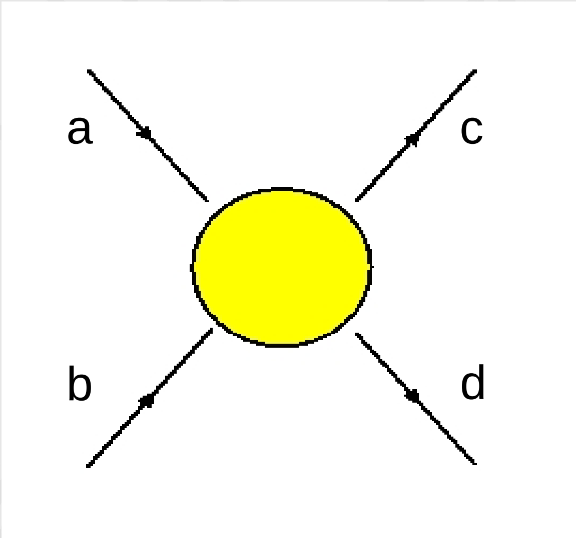
$$\hat{t} = (p_c - p_a)^2 = -\frac{\hat{s}}{2}(1-\cos\theta)$$

$$\hat{u} = (p_c - p_b)^2 = -\frac{\hat{s}}{2}(1+\cos\theta)$$

$$\hat{s} = (p_a + p_b)^2 = M_W^2$$

$$\cos\theta = \pm\frac{\sqrt{A^2-1}}{A}$$

$$A = M_W / 2E_T$$



W (Z) Diffractive cross sections

- $W^{+(-)}$ diffractive cross section

$$\frac{d\sigma}{d\eta_{e^-(e^+)}} = \sum_{a,b} \int dx_{IP} g(x_{IP}) \int dE_T \, f_{a/IP}(x_a) f_{b/\bar{p}}(x_b) \left[\frac{V_{ab}^2 G_F^2}{6s\Gamma_W M_W} \right] \frac{\hat{t}^2(\hat{u}^2)}{\sqrt{A^2 - 1}}$$

- Z^0 diffractive cross section

$$\sigma = \sum_{a,b} \int \frac{dx_{IP}}{x_{IP}} \int \frac{dx_b}{x_b} \int \frac{dx_a}{x_a} \bar{f}(x_{IP}) f_{a/IP}(x_a, \mu^2) f_{b/\bar{p}}(x_b, \mu^2) \left[\frac{2\pi C_{ab}^Z G_F M_Z^2}{3\sqrt{2}s} \right] \frac{d\hat{\sigma}(ab \to ZX)}{d\hat{t}}$$

- $f_{a/IP}$ is the quark distribution in the IP ➡ parametrization of the IP structure function (H1)

- $g(x_{IP})$ is the IP flux integrated over t

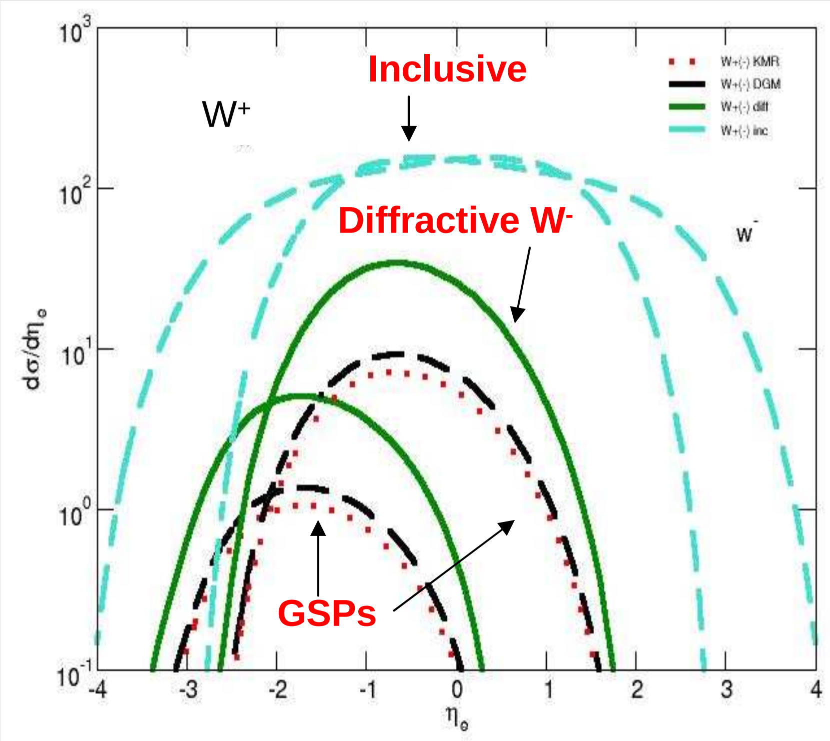
$$\bar{f}(x_{IP}) = \int_{-\infty}^{0} f_{IP/p}(x_{IP}, t) dt$$

$$C_{qq'}^Z 1/2 - 2|e_q|\sin^2\theta_W + 4|e_q|^2 \sin^4\theta_W$$

- θ_W is the Weinberg or weak-mixing angle

W^+ and W^- Cross Sections

Tevatron [sqrt (s) = 1.8 TeV]



IS + GSP models

	Pseudo-rapidity	Data (%)	R(%)	
1.8 TeV	$\vert\eta_e\vert<1.1$	1.15 ± 0.55	0.715 ± 0.045	CDF
		1.08 ± 0.25	0.715 ± 0.045	D0
	$1.5<\vert\eta_e\vert<2.5$	0.64 ± 0.24	1.7 ± 0.875	D0
	Total $W\to e\nu$	0.89 ± 0.25	0.735 ± 0.055	D0
	Total $Z\to e^+e^-$(*)	1.44 ± 0.80	0.71 ± 0.05	D0

GSP is an average of KMR (S^2 = 0.09) and GLM (S^2 = 0.086) estimations

$$R=\frac{\int_{-\eta}^{\eta}\sigma_{diff}^{W^+}+\sigma_{diff}^{W^-}}{\int_{-\eta}^{\eta}\sigma_{inc}^{W^+}+\sigma_{inc}^{W^-}}$$

- Ranges
  - $|\eta_e| < 1.1$
  - $1.5< |\eta_e|<2.5$

- **Tevatron, without GSP – 7.2 %**

* $|\eta|<1.1$

Quarkonium production in NRQCD

MBGD, M. M. Machado, M. V. T. Machado, PLB 683, 150-153 (2010)

Diffractive hadroproduction

o Focus on the following single diffractive processes

$$pp \to p + (J/\psi + \gamma) + X \qquad\qquad pp \to p + (Y + \gamma) + X$$

o Diffractive ratios as a function of transverse momentum p_T of quarkonium state

o Quarkonia produced with large p_T → **easy to detect**

o Singlet contribution

$$g + g \to \gamma + (c\bar{c})(^3S_1^{(1)})(\to J/\psi)$$

o Octet contributions

$$g + g \to \gamma + (c\bar{c})(^1S_0^{(8)})(\to J/\psi),$$

$$g + g \to \gamma + (c\bar{c})(^3P_J^{(8)})(\to J/\psi),$$

o Higher contribution on high p_T

J/ψ+γ production

- ✓ Considering the Non-relativistic Quantum Chromodynamics (NRQCD)
- ✓ Gluons fusion dominates over quarks annihilation
- ✓ Leading Order cross section ⟹ convolution of the partonic cross section with the PDF
- ✓ MRST 2001 LO ⟹ no relevant difference using MRST 2002 LO and MRST 2003 LO
- ✓ Non-perturbative aspects of quarkonium production
  - Expansion in powers of *v*
    - *v* is the relative velocity of the quarks in the quarkonia

NLO expansions in α_s one virtual correction and three real corrections

T. Mehen, Phys. Rev. D55 (1997) 4338

NRQCD Factorization

$$g + g \rightarrow \gamma + (c\bar{c}) \left[{}^3S_1^1, {}^3S_1^8\right]$$
$$g + g \rightarrow \gamma + (c\bar{c}) \left[{}^1S_0^8, {}^3P_J^8\right]$$

- Negligible contribution of quarks annihilation at high energies

$$\frac{d^2\sigma_{\rm inc}}{dy\,dp_T} = \int dx_1 g_p(x_1, \mu_F^2) g_p(x_2, \mu_F^2) \frac{4x_1 x_2 p_T}{2x_1 - \bar{x}_T e^y} \frac{d\hat{\sigma}}{d\hat{t}}$$

$$\bar{x}_T = 2m_T/\sqrt{s}$$

$$m_T = \sqrt{p_T^2 + m_\psi^2}$$

J/ψ rapidity

9.2 GeV^2

$\sqrt{s}$ is the center mass energy (LHC = 14 TeV)

NRQCD factorization

✓ x_1 (x_2) is the momentum fraction of the proton carried by the gluon

$M^2/s \leq x_1 < 1$ $\quad M$ → invariant mass of J/ψ+γ system

$$x_2 = \frac{x_1 \bar{x}_T e^{-y} - 2\tau}{2x_1 - \bar{x}_T e^{y}} \qquad \tau = \frac{m_\psi^2}{s}$$

✓ Cross section written as

$$\sigma(H) = \sum_n c_n \langle 0|O_n^H|0\rangle$$

Coefficients are computable in perturbation theory

Matrix elements of NRQCD operators

T. Mehen, Phys. Rev. D55 (1997) 4338

Matrix elements

$$\langle 0|O_n^H|0\rangle = \sum_X \sum_\lambda \langle 0|\kappa_n^\dagger|H(\lambda)+X\rangle\langle H(\lambda)+X|\kappa_n|0\rangle$$

Bilinear in heavy quarks fields which create as a pair $Q\overline{Q}$ — Quarkonium state

$$\frac{d\sigma}{dt}(g+g\to J/\psi+\gamma) = \frac{\pi^2 e_c^2 \alpha\alpha_s^2 m_c}{s^2}\left[\frac{10}{9}\left(\frac{s^2s_1^2+t^2t_1^2+u^2u_1^2}{s_1^2t_1^2u_1^2}\right)\langle O_8^{J/\psi}(^3S_1)\rangle + \frac{16}{27}\left(\frac{s^2s_1^2+t^2t_1^2+u^2u_1^2}{s_1^2t_1^2u_1^2}\right)\langle O_1^{J/\psi}(^3S_1)\rangle + \frac{3}{2}\frac{tu}{ss_1^2m_c^2}\langle O_8^{J/\psi}(^3S_1)\rangle + \frac{3}{2}\frac{1}{ss_1^2m_c^4}\left(2s(2m_c)^2+3tu-\frac{4tu(2m_c)^2}{s_1}\right)\langle O_8^{J/\psi}(^1P_0)\rangle\right],$$

$$s_1 = s-4m_c^2,\ t_1 = t-4m_c^2,\ u_1 = u-4m_c^2$$

$$e_c = \frac{2}{3}$$

α_s running

E. Braaten, S. Fleming, A. K. Leibovich, Phys. Rev. D63 (2001) 094006
F. Maltoni *et al.*, Phys. Lett. B638 (2006) 202

Matrix elements (GeV^3)

$\langle O_1^{J/\psi}(^3S_1)\rangle$	1.16	$\langle O_1^{\Upsilon}(^3S_1)\rangle$	10.9
$\langle O_8^{J/\psi}(^3S_1)\rangle$	1.19×10^{-2}	$\langle O_8^{\Upsilon}(^3S_1)\rangle$	0.02
$\langle O_8^{J/\psi}(^1S_0)\rangle$	0.01	$\langle O_8^{\Upsilon}(^1S_0)\rangle$	0.136
$\langle O_8^{J/\psi}(^1P_0)\rangle$	$0.01 \times m_c^2$	$\langle O_8^{\Upsilon}(^1P_0)\rangle$	0

$e_b = -\frac{1}{3}$

$m_b = 4.5$ GeV

$m_\Upsilon = 9.46$ GeV/c^2

Diffractive cross section

$$\frac{d^2\sigma_{SD}}{dy\,dp_T} = \int_{x_{\mathbb{P}}^{min}}^{x_{\mathbb{P}}^{max}} dx_{\mathbb{P}} \int_{\frac{M^2}{s x_{\mathbb{P}}}}^{1} dx_1 \int_{-1}^{0} dt\, f_{\mathbb{P}/p}(x_{\mathbb{P}}, t)$$

$$\times \quad g_{\mathbb{P}}(x_{\mathbb{P}}, \mu_F^2) g_p(x_2, \mu_F^2) \frac{4 x_1 x_{\mathbb{P}} x_2 p_T}{2 x_1 x_{\mathbb{P}} - \bar{x}_T e^y} \frac{d\hat{\sigma}}{d\hat{t}}$$

Momentum fraction carried by the Pomeron

Squared of the proton's four-momentum transfer

Pomeron flux factor

$$f_{\mathbb{P}/p}(x_{\mathbb{P}}, t) \propto x_{\mathbb{P}}^{1-2\alpha(t)} F^2(t)$$

$$\alpha(t) = \alpha_{\mathbb{P}}(0) + \alpha'_{\mathbb{P}} t$$ Pomeron trajectory

Variables to DDIS

Cuts for the integration over x_{IP}

$$x_{\mathbb{P}}^{min} \leq x_{\mathbb{P}} \leq 0.05$$

$$x_{\mathbb{P}}^{min} = \frac{\bar{x}_T e^y - 2\tau}{\bar{x}_T e^{-y} - 2}$$

Scales

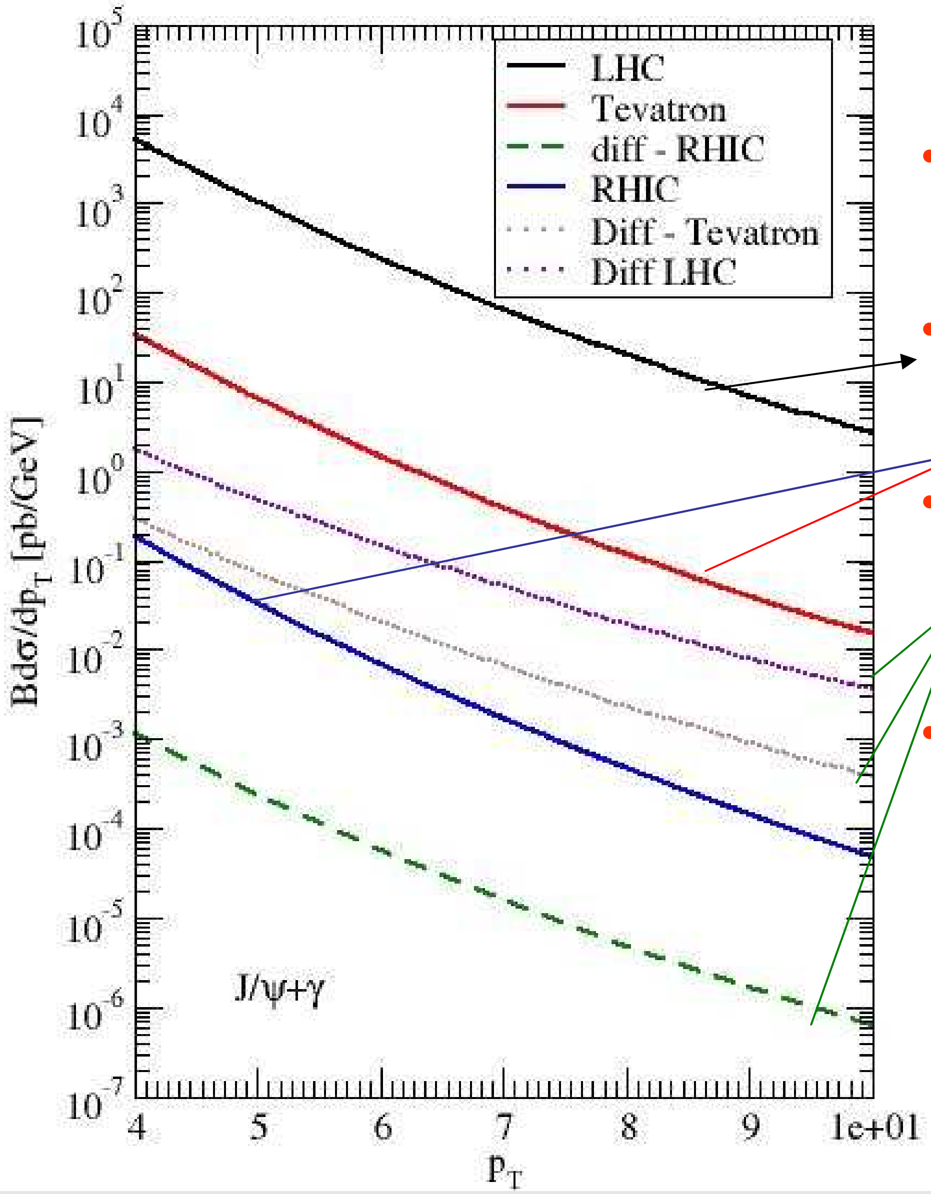
$$Q_0^2 = 2.5\, GeV^2 \qquad \Lambda_{QCD} = 0.2 \qquad \mu_F^2 = \frac{(p_T^2 + m_\psi^2)}{4}$$

$$x_2 = \frac{x_1 x_{\mathbb{P}} \bar{x}_T e^{-y} - 2\tau}{2x_1 x_{\mathbb{P}} - \bar{x}_T e^y},$$

$$\hat{s} = x_1 x_2 x_{\mathbb{P}} s, \qquad \hat{t} = m_\psi^2 - x_2 \sqrt{s} m_T e^y$$

$$\hat{u} = m_\psi^2 - x_1 x_{\mathbb{P}} \sqrt{s} m_T e^{-y}.$$

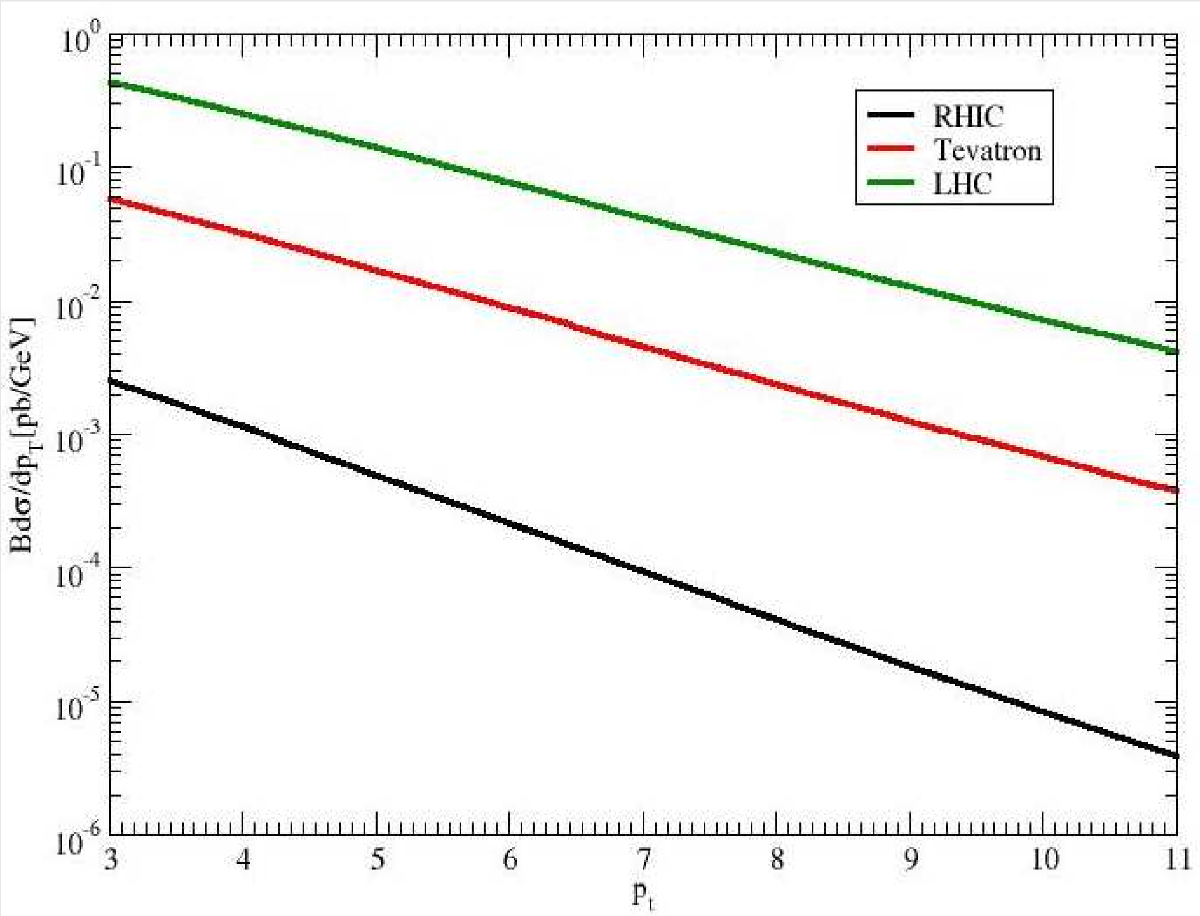
Results for J/ψ+γ



- Predictions for inclusive and diffractive cross sections
- LHC, Tevatron and RHIC
- Diffractive cross sections **considering GSP** ($\langle|S|^2\rangle$)
- B = 0.0594 is the branching ratio into electrons

$1 \leq p_T \leq 20$ at LHC

Results for Υ+γ

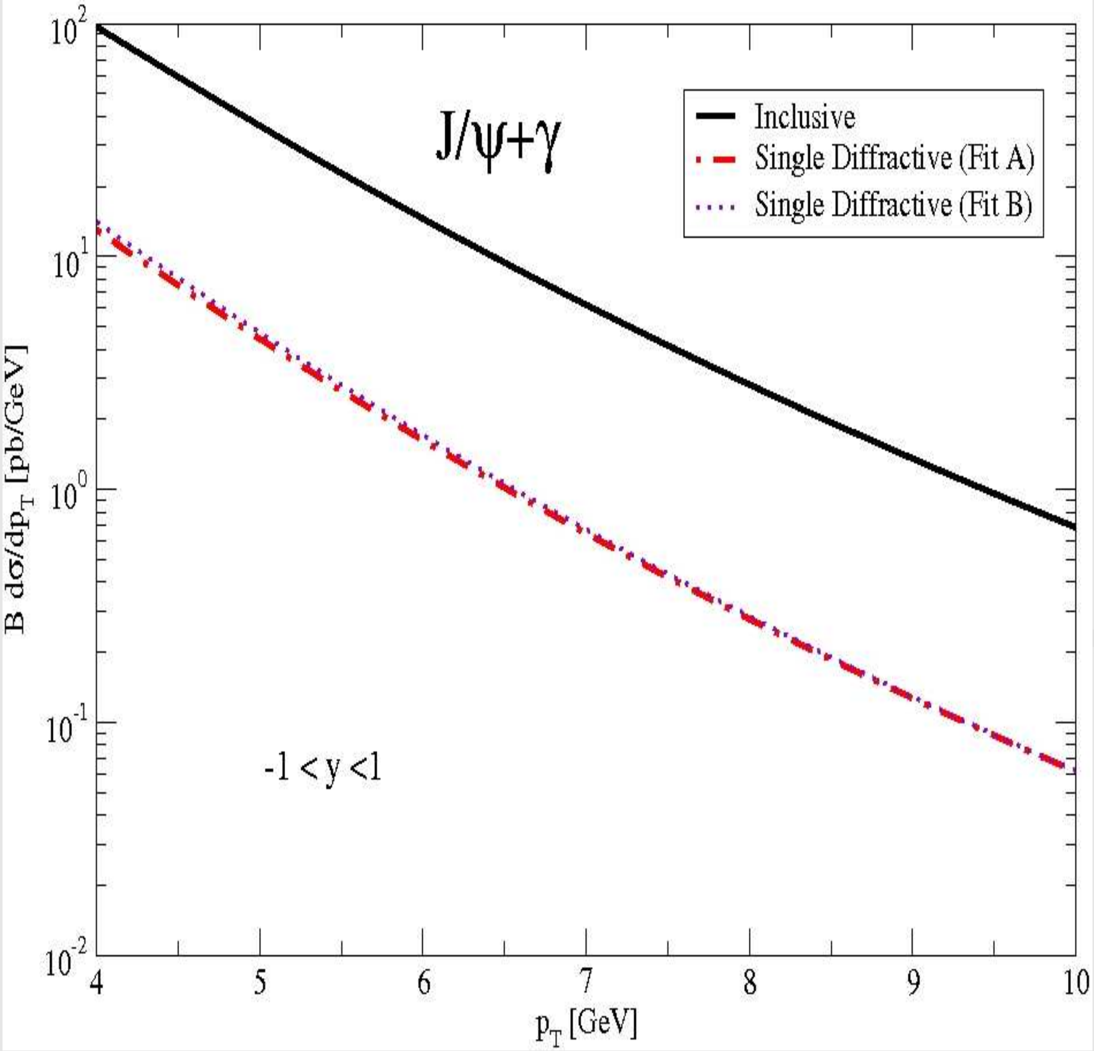


- Predictions of inclusive cross section
- **LHC**, Tevatron and RHIC

  $-1 < |y| < 1$

- B = 0.0238 is the branching ratio into electrons

Results for J/ψ+γ at LHC



- B = 0.0594
- Absolute value cross section strongly dependent

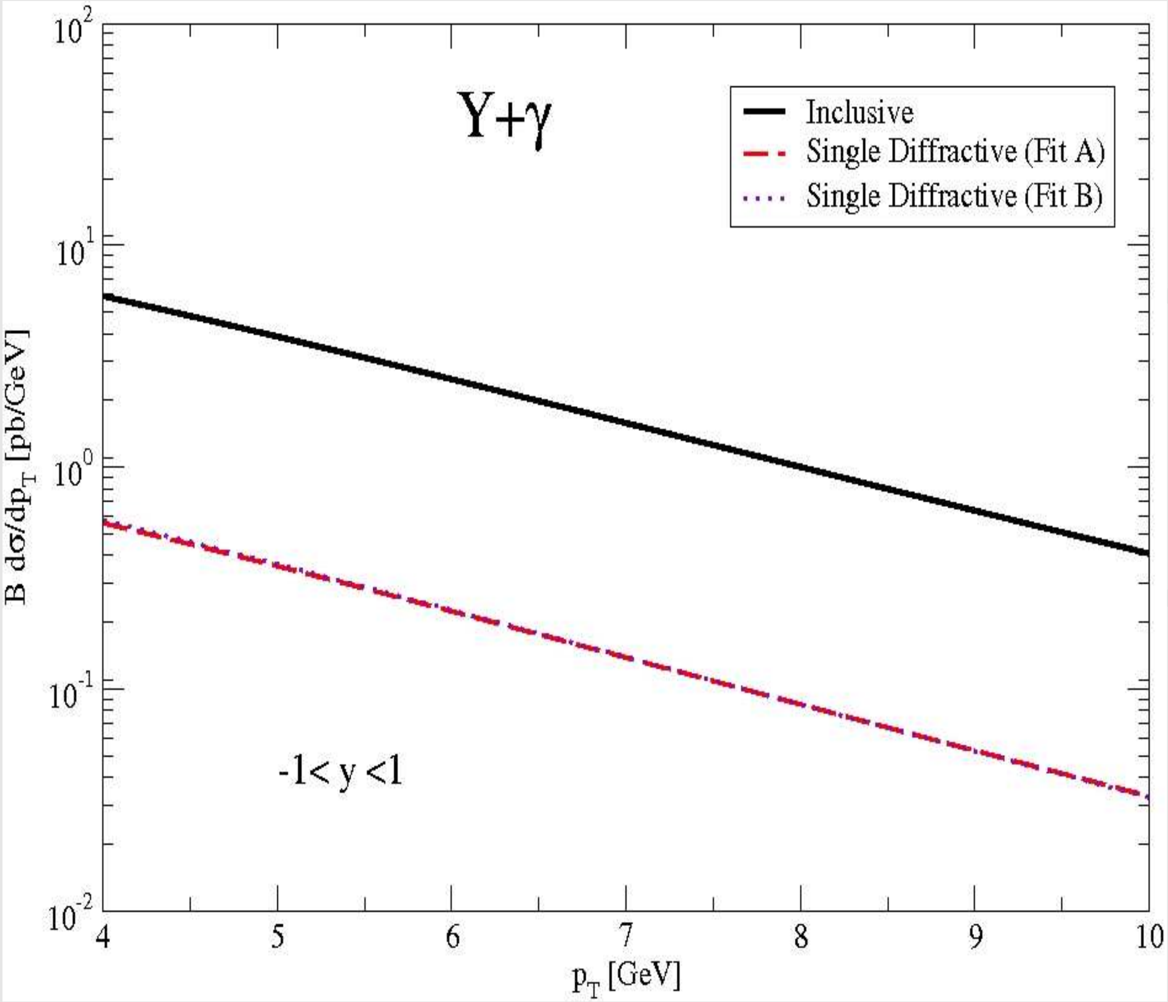
Quark mass

NRQCD matrix elements

Factorization scale

- Diffractive cross sections (DCS) without GSP (<|S|^2>)
- Comparison between two different sets of diffractive gluon distribution (H1)

Results for Υ+γ at LHC



- B = 0.0238
- Absolute value cross section strongly dependent

Quark mass

NRQCD matrix elements

Factorization scale

- Diffractive cross sections (DCS) without GSP ($<|S|^2>$)
- Comparison between two different sets of diffractive gluon distribution (H1)

Diffractive ratio at LHC

p_T [GeV]	4	6	8	10
$\frac{d\sigma_{\rm inc}}{dp_T}$ (J/Ψ)	97.04	14.54	2.82	0.68
$\frac{d\sigma_{\rm SD}}{dp_T}$ (J/Ψ)	0.78	0.10	0.017	0.0036
$R_{\rm SD}$ [%] (J/Ψ)	0.8	0.69	0.6	0.53
$\frac{d\sigma_{\rm inc}}{dp_T}$ (Υ)	5.91	2.49	1.00	0.41
$\frac{d\sigma_{\rm SD}}{dp_T}$ (Υ)	0.036	0.013	0.0054	0.0018
$R_{\rm SD}$ [%] (Υ)	0.6	0.53	0.54	0.44

** C. S. Kim, J. Lee and H. S. Song, Phys Rev D59 (1999) 014028

$J/\psi + \gamma$

P_T (GeV)	4	5	6	7	8	9	10
R $(P_T)(\%)$	0.52	0.52	0.50	0.48	0.47	0.46	0.44

- Slightly large diffractive ratio in comparison to **

❖ Could explain the p_T dependence in our results

[σ] = pb

considering FIT A

This work	Ref **
$\mu_F = \sqrt{\frac{(p_T^2 + m_\psi^2)}{4}}$	$\mu_F = E_T$
$<\|S\|^2>=0.06$	Renormalized Pomeron flux
Q^2 evolution in the gluon density	No Q^2 evolution in the gluon density

Heavy quark production

MBGD, M. M. Machado, M. V. T. Machado, PRD. 81, 054034 (2010)

MBGD, M. M. Machado, M. V. T. Machado, PRC. 83, 014903 (2011)

M. L. Mangano et al, Nucl. Phys. B 373, 295 (1992)

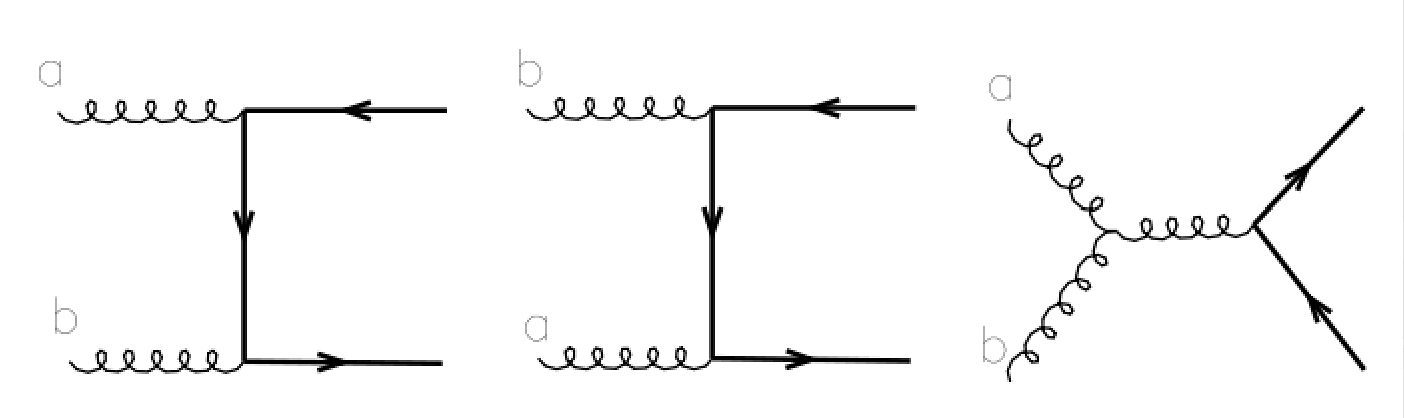
Heavy quark hadroproduction

o Focus on the following single diffractive processes

$$pp \rightarrow p + (c\bar{c}) + X \qquad pp \rightarrow p + (b\bar{b}) + X$$

o Diffractive ratios as a function of energy center-mass E_{CM}

$$g + g \rightarrow Q + \bar{Q}$$



o Diagrams contributing to the lowest order cross section

Total cross section LO

$$\sigma_{h_1h_2}(s,m_Q^2) = \sum_{i,j}\int_\rho^1 dx_1 \int_{\frac{\rho}{x_1}}^1 dx_2\, f_i^{h_1}(x_1,\mu_F^2) f_j^{h_2}(x_2,\mu_F^2)\hat\sigma_{ij}(\hat s, m_Q^2,\mu_F^2,\mu_R^2)$$

$$\rho = \frac{4m^2}{\hat s}$$

$$\hat s = x_1 x_2 s$$

$x_{1,2}$ are the momentum fraction

$f_i^{h_1}(x_1,\mu_F^2) f_j^{h_2}(x_2,\mu_F^2)$ are the parton distributions inner the hadron i=1 and j=2

Partonical cross section

$$\hat\sigma_{ij}(\hat s, m^2, \mu^2) = \frac{\alpha_S^2(\mu^2)}{m^2} f_{ij}\left(\rho, \frac{\mu^2}{m^2}\right)$$

$\mu_F(\mu_R)$ → factorization (renormalization) scale

$$\alpha_S = \frac{g^2}{4\pi}$$

$$\hat\sigma(gg \to Q\bar Q) = \sigma_0 \left(\frac{1}{NV}\right)\left[3\mathcal{L}(\beta)\xi_0 + 2(V-2)(1+\rho) + \rho(6\rho - N^2)\right]$$

$$\sigma_0 = \frac{\alpha_s^2}{m^2}\frac{\pi\beta}{24}\rho \qquad \mathcal{L}(\beta) = \frac{1}{\beta}log\left(\frac{1+\beta}{1-\beta}\right) - 2 \qquad \beta = \sqrt{1-\rho}$$

M. L. Mangano, P. Nason, G. Ridolfi Nucl. Phys. B373 (1992) 295

NLO Production

$$g+g \to Q+\overline{Q}+g$$

$$\hat{\sigma}_{ij}(\hat{s}, m_Q^2, \mu_F^2, \mu_R^2) = \frac{\alpha_s^2(\mu_R)}{m_Q^2} \sum_{k=0}^{\infty} [4\pi\alpha_s(\mu_R)]^k \sum_{l=0}^{k} f_{ij}^{(k,l)}(\rho) \ln^l\left(\frac{\mu_F^2}{m_Q^2}\right)$$

$$f_{gg}(\rho, \mu^2/m^2) = f_{gg}^{(0)}(\rho) + g^2(\mu^2)\left[f_{gg}^{(1)}(\rho) + \bar{f}_{gg}^{(1)}(\rho)\ln(\mu^2/m^2)\right] + O(g^4)$$

Running of the coupling constant

$$\frac{d\alpha_S(\mu^2)}{d\ln(\mu^2)} = -b_0\alpha_S^2 - b1\alpha_S^3 + O(\alpha_S^4)$$

$$b_0 = \frac{33 - 2n_{1f}}{12\pi}, \qquad b_1 = \frac{153 - 19n_{1f}}{24\pi^2}$$

n_{1f} = 3 (4) charm (bottom)

$$f_{gg}^{(0)}(\rho) = \frac{\pi\beta\rho}{192}\left[\frac{1}{\beta}(\rho^2 + 16\rho + 16)\ln\left(\frac{1+\beta}{1-\beta}\right) - 28 - 31\rho\right]$$

P. Nason, S. Dawson, R. K. Ellis Nucl. Phys. B303 (1988) 607

NLO functions

$$
\begin{aligned}
f_{gg}^{(1)} &= \frac{7}{1536\pi}\left[12\beta\ln^2(8\beta^2) - \frac{366}{7}\beta\ln(8\beta^2) + \frac{11}{42}\pi^2\right] + \beta\left[a_0 + \beta^2(a_1\ln(8\beta^2) + \right.\\
&+ \left. a_3\beta^4\ln(8\beta^2) + \rho^2(a_4\ln\rho + a_5\ln^2\rho) + \rho(a_6\ln\rho + a_7\ln^2\rho)\right]\\
&+ (n_{1f} - 4)\frac{\rho^2}{1024\pi}\left[\ln\left(\frac{1+\beta}{1-\beta}\right) - 2\beta\right]
\end{aligned}
$$

a_0	0.108068	a_4	0.0438768
a_1	-0.114997	a_5	-0.0760996
a_2	0.0428630	a_6	-0.165878
a_3	0.131429	a_7	-0.158246

Auxiliary functions

$$\beta = \sqrt{1-\rho}$$

$$
\begin{aligned}
\bar{f}_{gg}^{(1)}(\rho) &= \frac{1}{8\pi^2}\left[\left\{2\rho(59\rho^2 + 198\rho - 288)\ln\left(\frac{1+\beta}{1-\beta}\right) + 12\rho(\rho^2 + 16\rho + 16)h_2(\beta)\right.\right.\\
&- \left.\left. 6\rho(\rho^2 - 16\rho + 32)h_1(\beta) - \frac{4}{15}\beta(7449\rho^2 - 3328\rho + 724)\right\} + 12 f_{gg}^{(0)}(\rho)\ln\left(\frac{\rho}{4\beta^2}\right)\right]
\end{aligned}
$$

Diffractive cross section

$$\sigma^{\rm SD}_{h_1 h_2}(s, m_Q^2) = \sum_{i,j=q\bar{q},g} \int_\rho^1 dx_1 \int_{\rho/x_1}^1 dx_2$$
$$\times \int_{x_1}^{x_{I\!P}^{\max}} \frac{dx_{I\!P}^{(1)}}{x_{I\!P}^{(1)}} \bar{f}_{I\!P/h_1}\left(x_{I\!P}^{(1)}\right) f_{i/I\!P}\left(\frac{x_1}{x_{I\!P}^{(1)}}, \mu^2\right) f_{j/h_2}(x_2, \mu^2)\, \hat{\sigma}_{ij}(\hat{s}, m_Q^2, \mu^2) + (1 \rightleftharpoons 2)$$

$\bar{f}_{I\!P/h_1}\left(x_{I\!P}^{(1)}\right)$ → Pomeron flux factor

$$\beta = \frac{x}{x_{IP}}$$

$\beta f_{a/I\!P}(\beta, \mu^2)$ → Pomeron Structure Function (H1)

KKMR model → <|S|^2> = 0.06 at LHC single diffractive events

Parametrization of the pomeron flux factor and structure function → H1 Collaboration

Heavy quarks production at the LHC

Heavy Quark	$\sigma_{\text{inc}}(\sqrt{s} = 14\,\text{TeV})$	$\sigma_{\text{diff}}(\sqrt{s} = 14\,\text{TeV})$	R_{diff}
$c\bar{c}$	7811 $[\mu b]$	178 $[\mu b]$	2.3 %
$b\bar{b}$	393 $[\mu b]$	7 $[\mu b]$	1.7 %

Heavy quarks cross sections in NLO to pp collisions

GSP value decreases the diffractive ratio ($<|S|^2> = 0.06$)

Inclusive nuclear cross section at NLO

$$\sigma_A = A^2 \sigma_N$$

A_{PbPb} = 208 (5.5 TeV); 40 (6.3) TeV

$$\sigma^{\text{SD}}_{\text{pPb}} = 0.76\,(0.018)\,\text{mb}$$

charm (bottom)

$$\sigma^{\text{DPE}}_{\text{PbPb}} = 32.5\,(0.32)\,\mu\text{b}$$

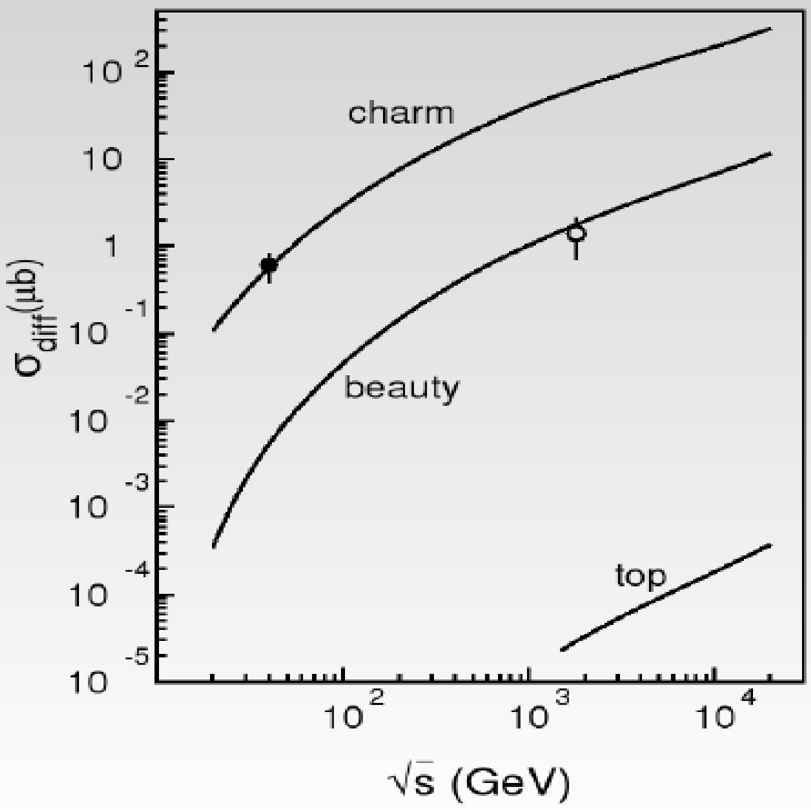
pA cross sections @ LHC

$\sigma^{\rm SD}_{\rm pPb} = 0.76\,(0.018)\,{\rm mb}$

charm (bottom)

$A_{\rm eff} = 4.39$

$S^2_{GAP} = 0.0287$



Similar results that

B. Kopeliovich *et al*, 0702106 [arXiv:hep-ph] (2007)

- Suppression factor

$$p + p \longrightarrow \bar{Q}QX + p$$

$$K = \left\{1 - \frac{1}{\pi}\,\frac{\sigma^{pp}_{tot}(s)}{B_{sd}(s) + 2B^{pp}_{el}(s)} + \frac{1}{(4\pi)^2}\,\frac{[\sigma^{pp}_{tot}(s)]^2}{B^{pp}_{el}(s)\,[B_{sd}(s) + B^{pp}_{el}(s)]}\right\}$$

than is suggested by Eq. (70). Therefore, the predicted energy dependence of the survival probability Eq. (70) might be quite wrong and the diffractive cross section at the LHC energy may be overestimated.

$$\sigma^{pA}_{sd}(pA \to X\bar{Q}QA) = A_{eff}\,\sigma^{pp}_{sd}(pA \to X\bar{Q}Qp)$$

σ_{pA} ~ 0.8 mb (charm)

$A_{eff} \approx 10.$

Diffractive cross sections @ LHC

Inclusive cross section

Nucleus-Nucleus collision

	Pb-Pb $(C\bar{C})$	Pb-Pb $(B\bar{B})$
$\sigma_A[mb]$	$188165, 16$	$7340, 23$

$A_{Pb} = 240$

Diffractive cross sections

Coherent	PbPb $(c\bar{c})$	PbPb $(b\bar{b})$
$\sigma_{\rm coh}/A^2$	3.7 mb	0.06 mb
$\sigma_{\rm coh}^{\rm abs}$	$9686 - 0.16$ mb	$156 - 0.003$ mb
$R_{\rm coh}\,[\%]$	86	35
$R_{\rm coh}^{\rm abs}\,[\%]$	$5.2 - 8.6 \times 10^{-5}$	$2.1 - 3.5 \times 10^{-5}$

- **Coherent**

Pomeron emmited by the nucleus

$$A + A \rightarrow X + A + [LRG] + A$$

$$F(t) \approx exp(R_A^2 t/6)$$

$R_A = r_0 A^{1/3}$ $r_0 = 1,2$ fm

- Predictions to cross sections possible to be verified at the LHC

Very small diffractive ratio

Diffractive cross sections @ LHC

Incoherent	PbPb ($c\bar{c}$)	PbPb ($b\bar{b}$)
$\sigma_{\rm inc}/A^2$	1.68 mb	0.03 mb
$\sigma_{\rm inc}^{\rm abs}$	$4356 - 0.07$ mb	$85 - 0.001$ mb
$R_{\rm inc}[\%]$	38	19
$R_{\rm inc}^{\rm abs}[\%]$	$2.28 - 3.8 \times 10^{-5}$	$1.14 - 1.9 \times 10^{-5}$

❖ **Incoherent**

Pomeron emmited by

a nucleon inner the nucleus

$$A + A \rightarrow X + A + [LRG] + A^*$$

$$\sigma_{A_{diff}} \approx A^2 \sigma_{N_{diff}}$$

❖ No values to <|S|^2> for single diffractive events in AA collisions

❖ Estimations to central Higgs production → <|S|^2> ~ 8 x 10^{-7}

❖ Values of diffractive cross sections possible to be verified experimentally

A_{Pb} = 240

DPE results at LHC

$Q\bar{Q}$	$\sigma_{\rm inc}$ [μb]	$\sigma_{\rm DPE}$ [μb]	$R_{\rm DPE}$ [%]
$c\bar{c}$	7811	13.6–0.53	0.17–7×10^{-3}
$b\bar{b}$	393	0.053–0.027	0.01–0.007

Ingelman-Schlein

Bialas-Landshoff

pp collisions LHC (14 TeV)

$S^2_{\rm gap} = 0.026$

	CaCa [$c\bar{c}$]	PbPb [$c\bar{c}$]	CaCa [$b\bar{b}$]	PbPb [$b\bar{b}$]
$\sigma_{AA}^{\rm DPE}$ [μb]	22.8–2.8	31.1–4.2	0.25–0.14	0.32–0.2
$R_{\rm coh}^{\rm DPE}$[%]	$3-0.4\times10^{-4}$	$2-0.2\times10^{-4}$	$8-4\times10^{-5}$	$4-3\times10^{-6}$

$S^2_{\rm gap} = 0.032\ (0.031)$ $A^2_{\rm eff} = 9.52\ (6.21)$

$$S^2_{\rm gap} = a/[b + \ln(\sqrt{s/s_0})]$$

$a = 0.126$ $b = -4.688$ $s_0 = 1\ {\rm GeV}^2$

Ingelman-Schlein > Bialas-Landshoff

AA collisions LHC

CaCa (6.3 TeV)

PbPb (5.5 Tev)

Higgs production

MBGD, M. M. Machado, G. G. Silveira, PRD. 83, 074005 (2011)

Higgs production

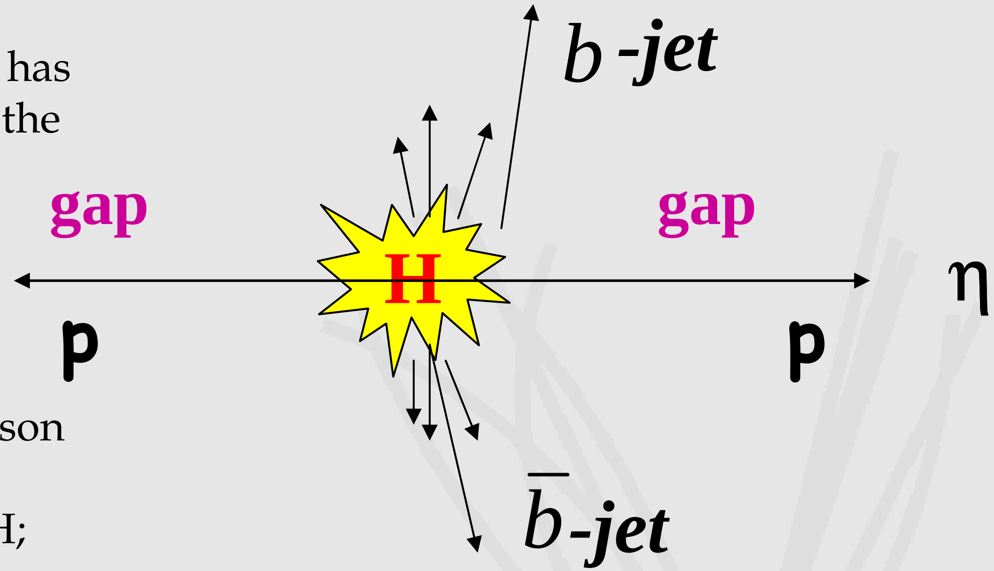
✓ Standard Model (SM) of Particle Physics has unified the Eletromagnetic interaction and the weak interaction;

✓ Particles acquire mass through their interaction with the Higgs Field;

✓ Existence of a new particle: the Higgs boson

✓The theory does not predict the mass of H;

✓ Predicts its production rate and decay modes for each possible mass;



- Exclusive diffractive Higgs production pp → p H p : **3-10 fb**
- Inclusive diffractive Higgs production p p → p + X + H + Y + p : **50-200 fb**

Albert de Roeck X BARIONS (2004)

Tevatron cuts

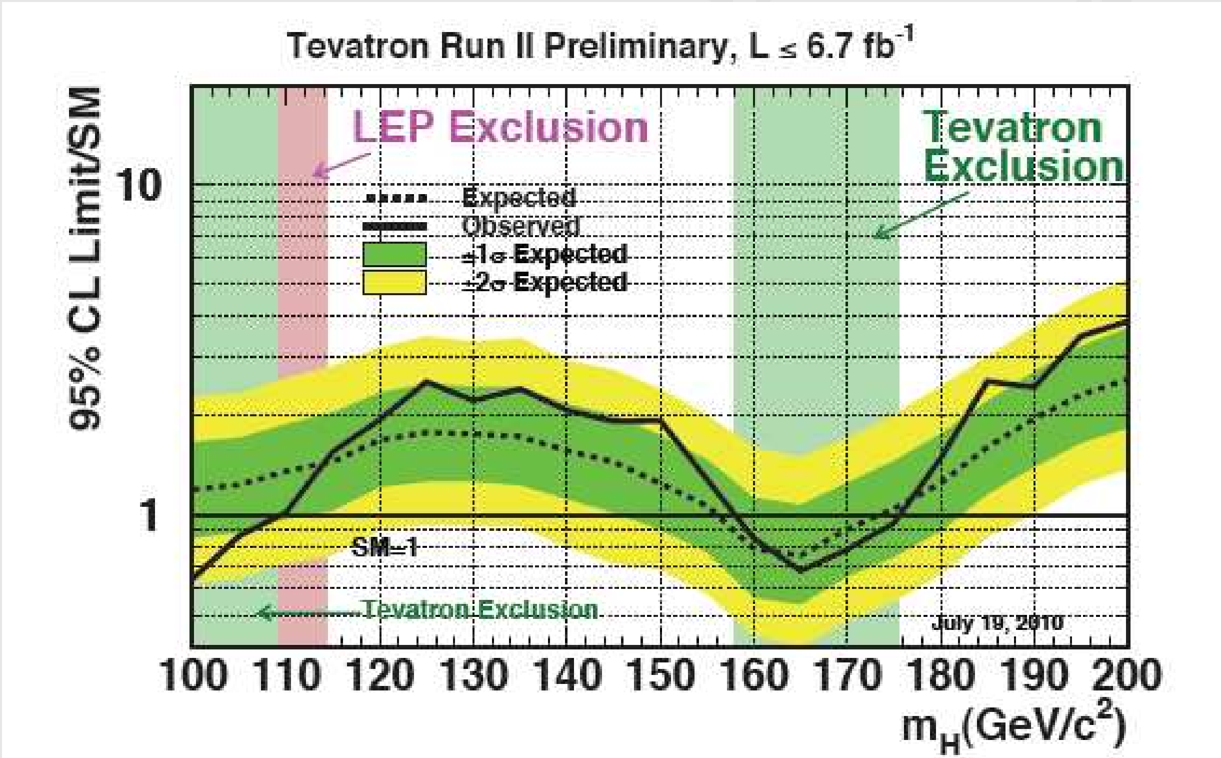
✓ LHC opens a new kinematical region:

✓ CM Energy in pp Collisions: 14 TeV   7x Tevatron Energy

✓ Luminosity: 10 – 100 fb^{-1} 10 x Tevatron luminosity

✓Evidences show new allowed mass range excluded for Higgs Boson production

✓Tevatron exclusion ranges are a combination of the data from CDF and D0



The TEVNPH Working Group, 1007.4587 [hep-ph]

D. Graudenz *et al.* PRL 70 (1993) 1372

Gluon fusion

o Focus on the gluon fusion

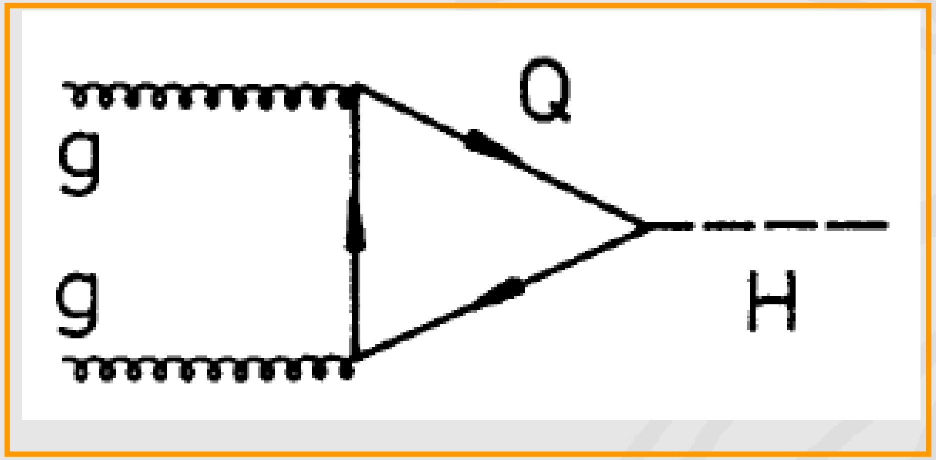
$$pp \longrightarrow gg \longrightarrow H$$

o Main production mechanism of Higgs boson in high-energy pp collisions

Gluon coupling to the Higgs boson in SM

triangular loops of top quarks

Lowest order to *gg* contribution



M. Spira *et al.* 9504378 [hep-ph]

Gluon fusion

Lowest order

partonic cross section expressed by the gluonic width of the Higgs boson

$$\hat{\sigma}_{LO}(gg \to H) = \frac{\sigma_0}{m_H^2}\delta(\hat{s} - m_H^2)$$

$$\sigma_0 = \frac{8\pi^2}{m_H^3}\Gamma_{LO}(H \to gg)$$

$$\Gamma_{LO}(H \to gg) = \frac{G_F \alpha_s^2}{36\sqrt{2}\pi^3} m_H^3 \left| \frac{3}{4} \sum_Q A_Q(\tau_Q) \right|^2$$

$$A_Q(\tau_Q) = 2[\tau + (\tau - 1)f(\tau)]/\tau^2$$

$$f(\tau) = \arcsin^2 \sqrt{\tau}$$

Quark Top

$\hat{s}$ → *gg* invariant energy squared

✓ dependence τ_Q → $\tau_Q = m_H^2/4m_Q^2$

LO hadroproduction

✓ Lowest order ➡ two-gluon decay width of the Higgs boson

$$\sigma_0 = \frac{G_F \alpha_s^2(\mu^2)}{288\sqrt{2}\pi} \left| \frac{3}{4} \sum_q A_Q(\tau_Q) \right|^2$$

✓Gluon luminosity ➡ $\frac{d\mathcal{L}^{gg}}{d\tau} = \int_\tau^1 \frac{dx}{x} g(x, M^2) g(\tau/x, M^2)$

PDFs
MSTW2008

✓ Lowest order proton-proton cross section

$$\sigma_{LO}(pp \to H) = \sigma_0 \tau_H \frac{d\mathcal{L}^{gg}}{d\tau_H}$$

✓ Renormalization scale μ_Q $\tau = \tau_H$ $\tau_H = \frac{m_H^2}{s}$

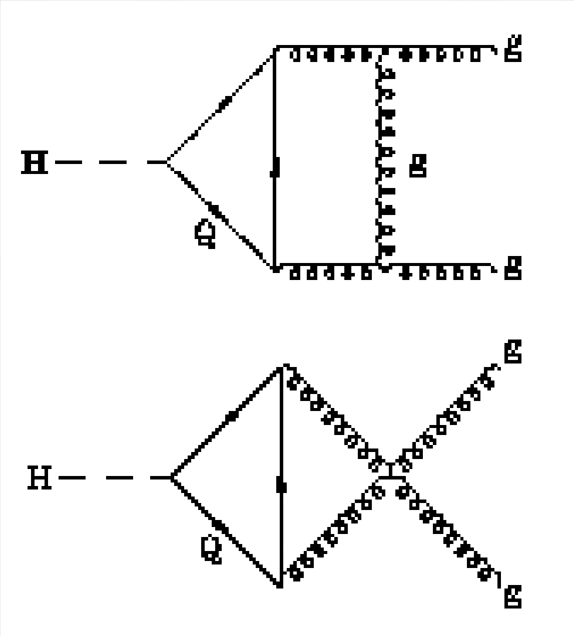
✓s ➡ invariant *pp* collider energy squared

Virtual diagrams

- Coefficient $C(\tau_Q)$ → contributions from the virtual two-loop corrections

- Regularized by the infrared singular part of the cross section for real gluon emission

$$C(\tau_Q) = \pi^2 + c(\tau_Q) + \frac{33 - 2N_F}{6} \ln \frac{\mu^2}{m_H^2}$$



✓ Infrared part

✓ Finite τ_Q dependent piece

✓ Logarithmic term depending on the renormalization scale μ

Delta functions

o Contributions from gluon radiation in *gg*, *gq* and *qq* scattering

o Dependence of the parton densities — renormalization scale μ, factorization scale *M*

$$\Delta\sigma_{gg} = \int_{\tau_H}^{1} d\tau \frac{d\mathcal{L}^{gg}}{d\tau} \times \frac{\alpha_s}{\pi}\sigma_0 \left\{ -\hat{\tau} P_{gg}(\hat{\tau}) \log\frac{M^2}{\hat{s}} + d_{gg}(\hat{\tau}, \tau_Q) + 12\left[\left(\frac{\log(1-\hat{\tau})}{1-\hat{\tau}}\right)_+ - \hat{\tau}[2-\hat{\tau}(1-\hat{\tau})]\log(1-\hat{\tau})\right]\right\}$$

$$\Delta\sigma_{gq} = \int_{\tau_H}^{1} d\tau \sum_{q,\bar{q}} \frac{d\mathcal{L}^{gq}}{d\tau} \times \frac{\alpha_s}{\pi}\sigma_0 \left\{ \hat{\tau} P_{gq}(\hat{\tau}) \left[-\frac{1}{2}\log\frac{M^2}{\hat{s}} + \log(1-\hat{\tau})\right] + d_{gq}(\hat{\tau}, \tau_Q)\right\}$$

$$\Delta\sigma_{q\bar{q}} = \int_{\tau_H}^{1} d\tau \sum_{q} \frac{d\mathcal{L}^{q\bar{q}}}{d\tau} \times \frac{\alpha_s}{\pi}\sigma_0 \, d_{q\bar{q}}(\hat{\tau}, \tau_Q)$$

$$\hat{\tau} = \tau_H/\tau$$

Renormalization scale

QCD coupling $\alpha_s(\mu^2)$ in the radiative corrections and LO cross sections

d functions

$$P_{gg}(\hat{\tau}) = 6\left\{\left(\frac{1}{1-\hat{\tau}}\right)_+ + \frac{1}{\hat{\tau}} - 2 + \hat{\tau}(1-\hat{\tau})\right\} + \frac{33-2N_F}{6}\delta(1-\hat{\tau})$$

$$P_{gq}(\hat{\tau}) = \frac{4}{3}\frac{1+(1-\hat{\tau})^2}{\hat{\tau}}$$

F_+ : usual + distribution

$$F(\hat{\tau})_+ = F(\hat{\tau}) - \delta(1-\hat{\tau})\int_0^1 d\hat{\tau}' F(\hat{\tau}')$$

$$\tau_Q = m_H^2/4m_Q^2 \ll 1$$

Considering only the heavy-quark limit

Region allowed by Tevatron combination

$$c(\tau_Q) \rightarrow \frac{11}{2}$$

$$d_{gg}(\hat{\tau}, \tau_Q) \rightarrow -\frac{11}{2}(1-\hat{\tau})^3$$

$$d_{gq}(\hat{\tau}, \tau_Q) \rightarrow -1 + 2\hat{\tau} - \frac{\hat{\tau}^2}{3}$$

$$d_{q\bar{q}}(\hat{\tau}, \tau_Q) \rightarrow \frac{32}{27}(1-\hat{\tau})^3$$

NLO Cross Section

❖ Gluon radiation ➡ two parton final states $gg \to H$

❖ Invariant energy $\hat{s} \geq m_H^2$ in the $gg, gq \text{ and } q\bar{q}$ channels

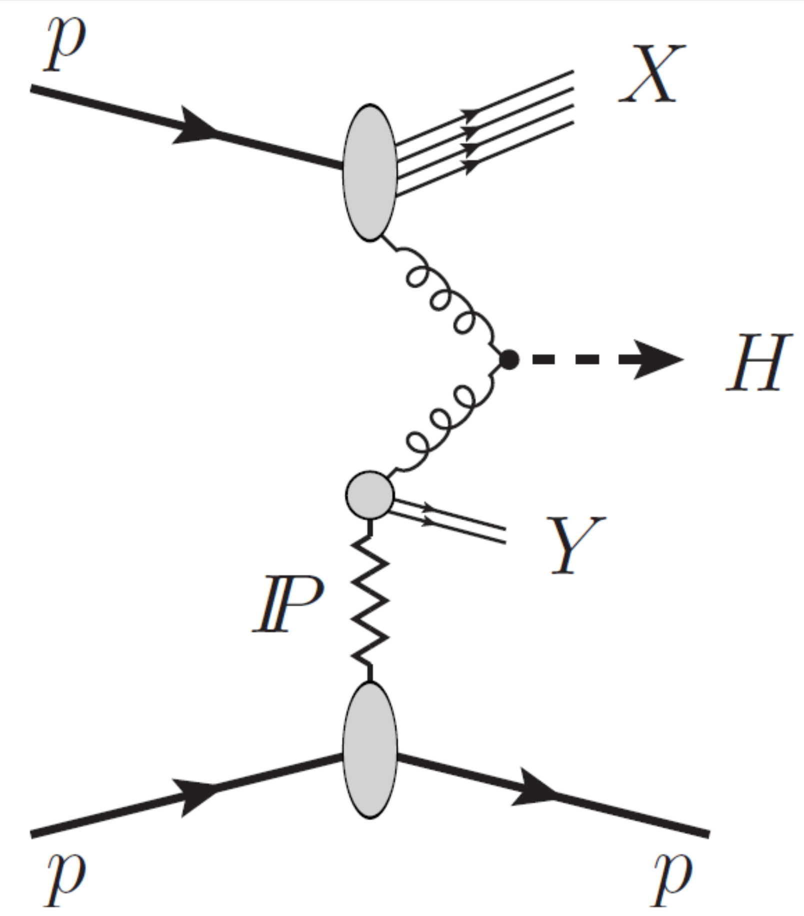
❖ New scaling variable $\hat{\tau}$ ➡ supplementing τ_H and τ_Q $\hat{\tau} = \frac{m_H^2}{\hat{s}}$

❖ The final result for the pp cross section at NLO

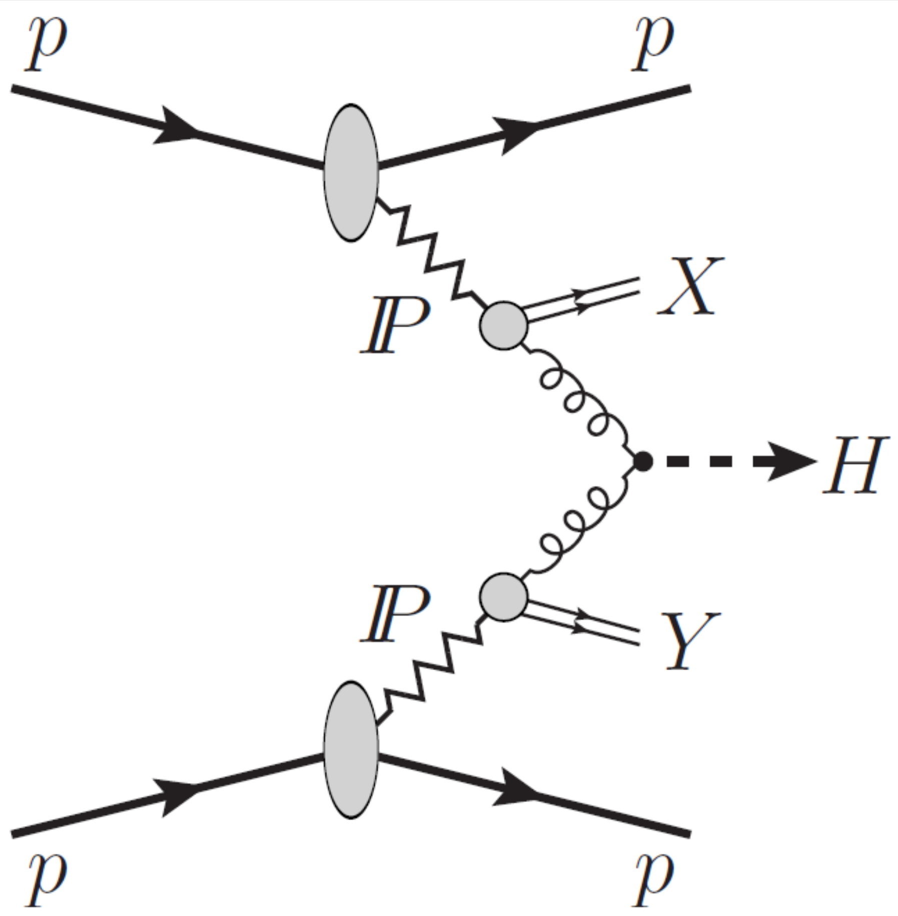
$$\sigma(pp \to H + X) = \sigma_0 \left[1 + C\frac{\alpha_s}{\pi}\right] \tau_H \frac{d\mathcal{L}^{gg}}{d\tau_H} + \Delta\sigma_{gg} + \Delta\sigma_{gq} + \Delta\sigma_{q\bar{q}}$$

❖ Renormalization scale in α_s and the factorization scale of the parton densities to be fixed properly

Diffractive processes



Single diffractive



Double Pomeron Exchange

Diffractive cross sections

Single diffractive

$$\sigma_{I\!Pp\to H}(M_H, M_X) = C_g \int_0^1 \int_0^1 F_{g/p}(\xi_p) \cdot F_{g/I\!P}(\beta) \cdot \sigma_{gg\to H}(M_H, \hat{s}) \; d\beta \, d\xi_p$$

Double Pomeron Exchange

$$\sigma_{I\!Pp\to H}(M_H, M_X) = C_g \int_0^1 \int_0^1 F_{g/I\!P_A}(\beta) \, F_{g/p_B}(\xi_p) \, \sigma_{gg\to H}(M_H, \hat{s}) \; d\beta \, d\xi_p$$

C_g → Normalization

Momentum fractions: pomeron and quarks

$$\xi = 1 - x_p \qquad \beta = \frac{x}{x_{IP}}$$

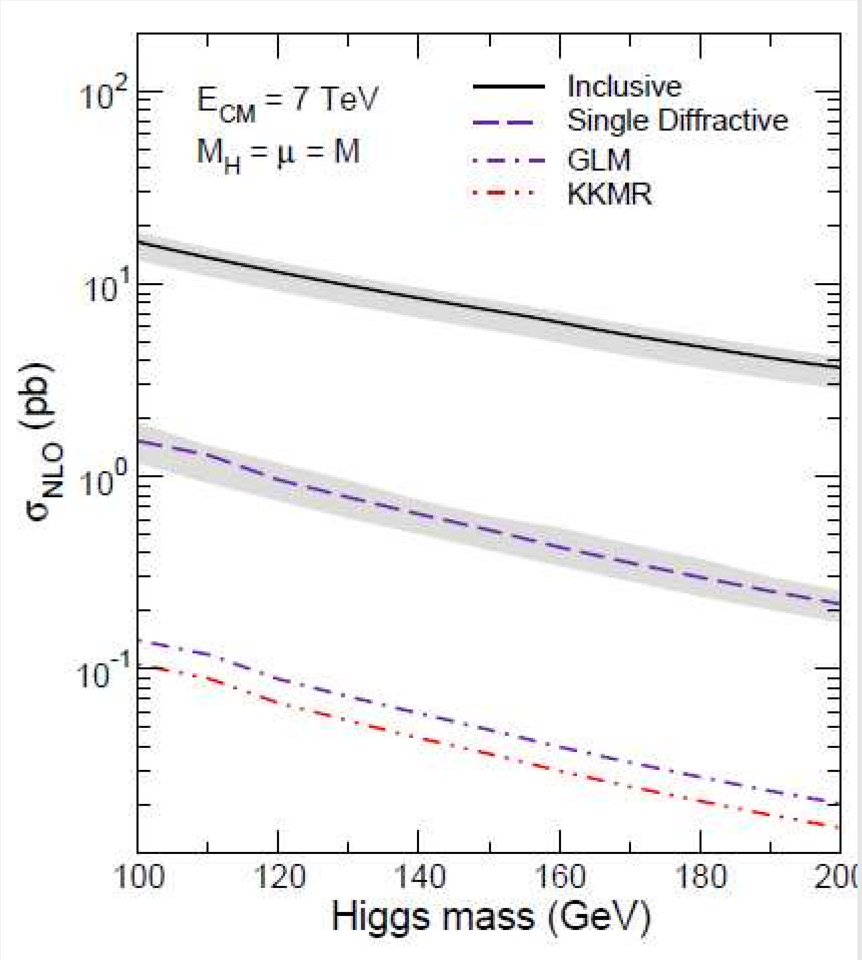
$F_{g/p}(\xi_p)$ → Gluon distributions in the proton — MSTW (2008)

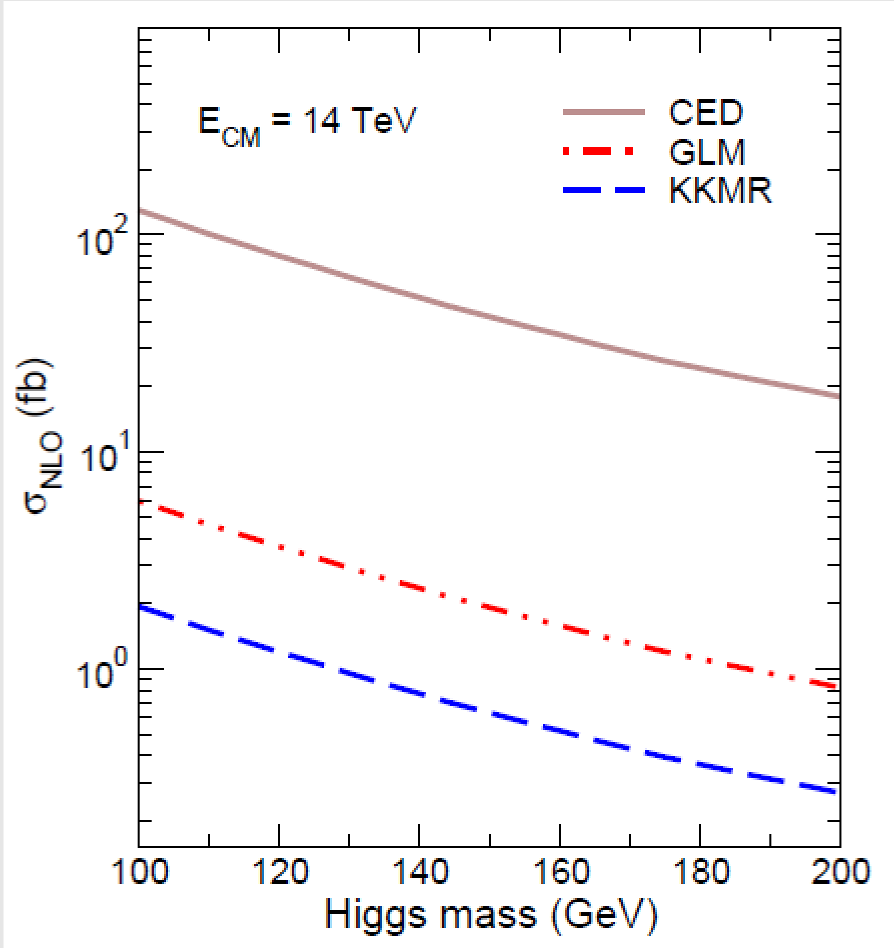
$F_{g/I\!P}(\beta)$ → $f_{I\!P/h}(x_{I\!P}) \, f_{i/I\!P}\left(\frac{x}{x_{I\!P}}, \mu^2\right)$ — H1 parametrization (2006)

Pomeron flux

Gluon distributions (i) in the Pomeron IP

FIT Comparison :: **SD** vs. **DPE**





SD production as M_H function (NLO)

Mass (GeV)	$\sqrt{s}$ (TeV)			
	1.96	7.	8.	14.
120	5.36(4.23)	88.59(66.44)	119.70(90.11)	346.43(256.62)
140	2.57(2.02)	58.69(44.02)	81.43(61.30)	248.75(184.26)
160	1.24(0.98)	39.56(29.67)	56.07(42.21)	183.06(135.60)
180	0.60(0.47)	27.60(20.70)	40.23(30.28)	134.46(99.60)
200	0.31(0.24)	19.96(14.97)	29.10(21.90)	104.65(77.52)

GLM KKMR

Exclusive Higgs boson production

MBGD, G. G. Silveira, Phys. Rev. D 78, 113005 (2008)
MBGD, G. G. Silveira, Phys. Rev. D 82, 073004 (2011)

Khoze, Martin, Ryskin, EPJB 401, 330 (1997)

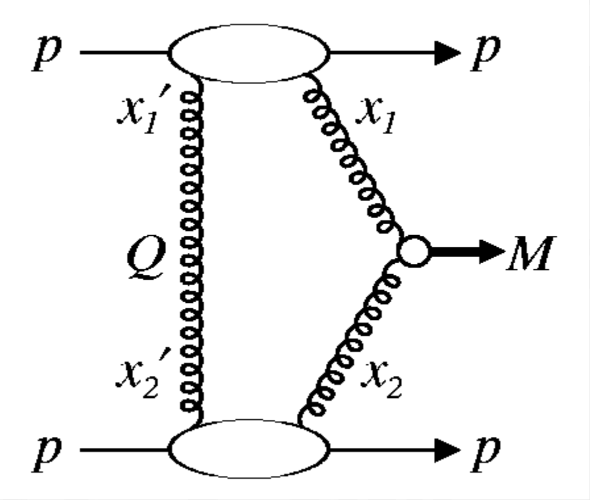
Diffractive Higgs Production

- The reaction $pp \rightarrow p + H + p$
-
- Protons lose small fraction of their energy :: **scattering in small angles**
- Nevertheless enough to produce the Higgs Boson

Durham Model

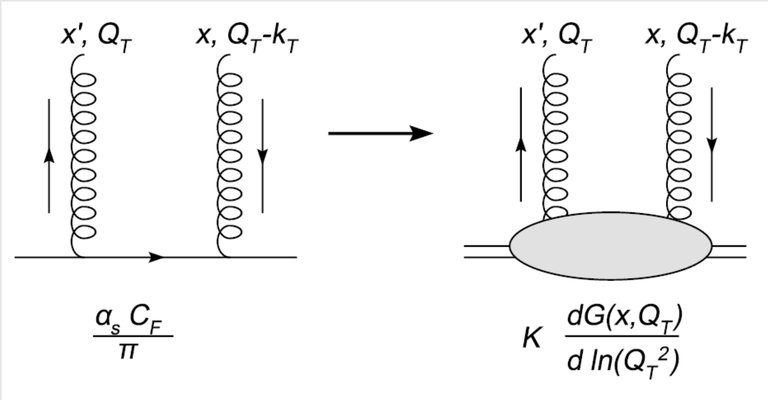
$$\frac{d\sigma}{dy} = \frac{|M|^2}{16^2 \pi^3 b^2}$$

G_F is the Fermi constant and $Q_T^2 \equiv -\mathbf{Q}_T^2$



Neglected the exchanged transverse momentum in the integrand

2-gluon emission



- The probability for a quark emit 2 gluon in the t-channel is given by the integrated gluon distribution

$$f(x,Q) \equiv K \partial G(x,Q) / \partial \ln Q^2$$

- The factor K is related to the non-diagonality of the distribution

$$K \approx e^{-bk_T^2/2} \frac{2^{2\lambda+3}}{\sqrt{\pi}} \frac{\Gamma(\lambda+5/2)}{\Gamma(\lambda+4)}$$

$$\frac{d\sigma}{dy} \approx \frac{\alpha_s G_F \sqrt{2}}{9b^2} \left[\int \frac{d^2 Q_T}{Q_T^4} f(x_1, Q_T) f(x_2, Q_T) \right]^2$$

Sudakov form fators

- The former cross section is **infrared divergent**!
- The regulation of the amplitude can be done by suppression of gluon emissions from the production vertex;
- The Sudakov form factors accounts for the probability of emission of one gluon

$$\frac{C_A \alpha_s}{\pi} \int_{Q_T^2}^{m_H^2/4} \frac{dp_T^2}{p_T^2} \int_{p_T}^{m_H/2} \frac{dE}{E} \sim \frac{C_A \alpha_s}{4\pi} \ln^2 \left(\frac{m_H^2}{Q_T^2} \right)$$

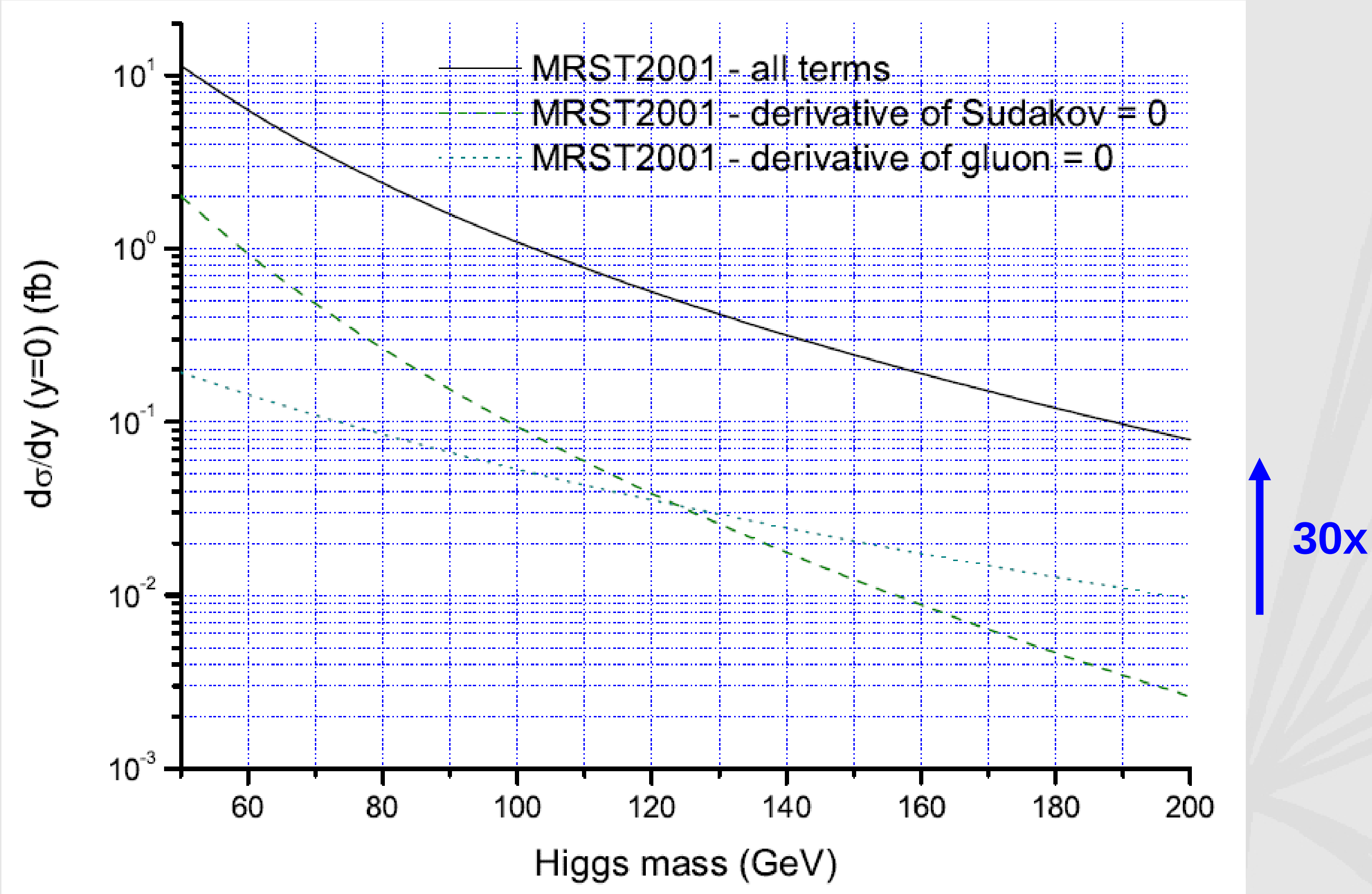
- The **suppression** of several gluon emissions exponentiate

$$e^{-S} = \exp\left(- \int_{Q_T^2}^{m_H^2/4} \frac{dp_T^2}{p_T^2} \frac{\alpha_s(p_T^2)}{2\pi} \int_0^{1-\Delta} dz \left[z P_{gg}(z) + \sum_q P_{qg}(z) \right] \right)$$

- Then, the gluon distributions are modified in order to include S

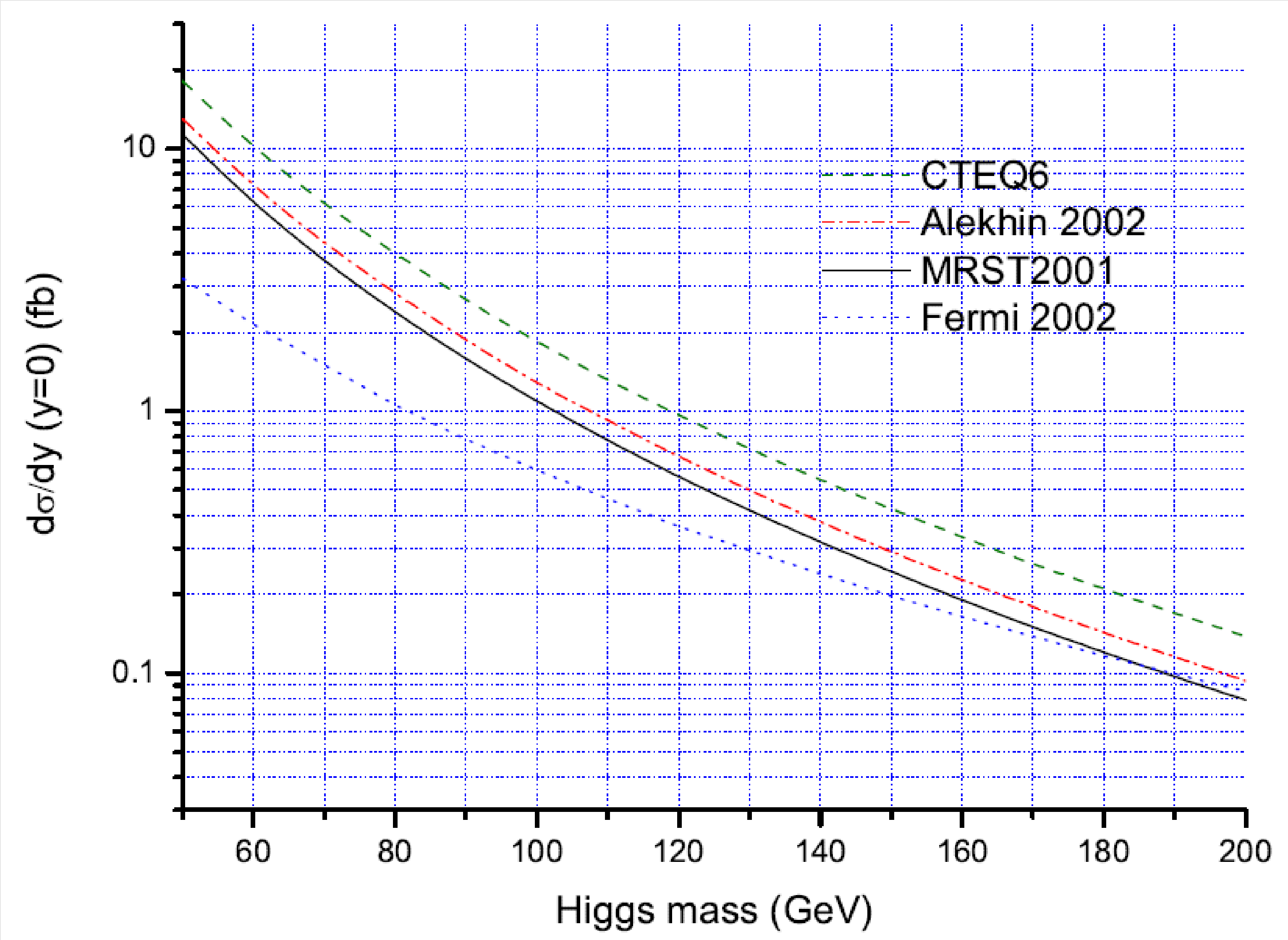
$$\tilde{f}(x, Q_T) = \frac{\partial}{\partial \ln Q_T^2} \left(e^{-S/2} \, G(x, Q_T) \right)$$

Cross section I :: Sudakov



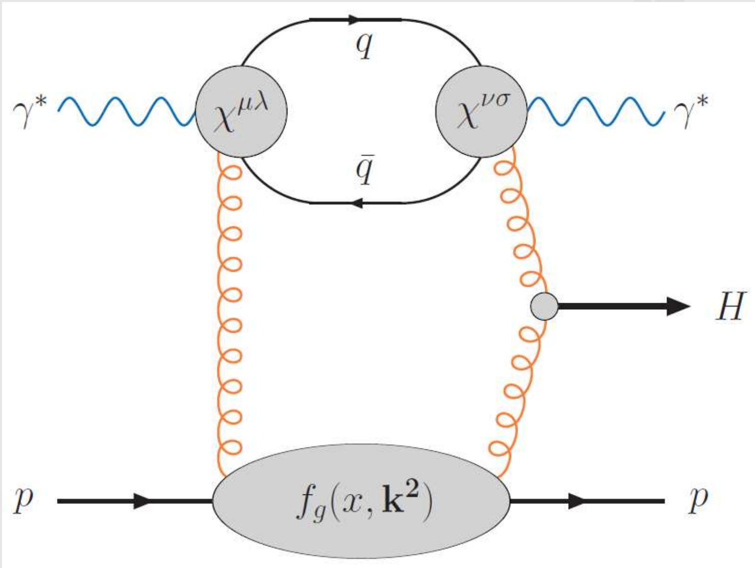
$$\frac{d\sigma}{dy} \approx \frac{\alpha_s G_F \sqrt{2}}{9b^2}\left[\int \frac{d^2 Q_T}{Q_T^4} \tilde{f}(x_1, Q_T)\tilde{f}(x_2, Q_T)\right]^2$$

Cross section II :: PDFs



Photoproduction mechanism

- The Durham group's approach is applied to the photon-proton process;
- This is a subprocess of **Ultraperipheral Collisions**;
- Hard process: photon splitting into a color dipole, which interacts with the proton;



Dipole contribution

$$\Im A_T = -\frac{s}{3}\frac{M_H^2\alpha_s^3\,\alpha}{\pi v}\sum_q e_q^2\left(\frac{2C_F}{N_c}\right)\int\frac{\mathrm{d}\boldsymbol{k}^2}{\boldsymbol{k}^6}\int_0^1\frac{[\tau^2+(1-\tau)^2][\alpha_\ell^2+(1-\alpha_\ell)^2]\boldsymbol{k}^2}{\boldsymbol{k}^2\tau(1-\tau)+Q^2\alpha_\ell(1-\alpha_\ell)}\,\mathrm{d}\alpha_\ell\,\mathrm{d}\tau.$$

γp cross section

- The cross section is calculated for central rapidity ($y_H = 0$)

$$\left.\frac{d\sigma}{dy_H\, dt}\right|_{y_H,t=0} = \frac{S^2_{gap}}{2\pi B}\left(\frac{\alpha_s^2 \alpha M_H^2}{3N_c \pi v}\right)^2 \left(\sum_q e_q^2\right)^2 \left[\int_{\mathbf{k}_0^2}^{\infty} \frac{d\mathbf{k}^2}{\mathbf{k}^6}\, e^{-S(\mathbf{k}^2, M_H^2)}\, f_g(x, \mathbf{k}^2)\, \mathcal{X}(\mathbf{k}^2, Q^2)\right]^2$$

- Proton content[1]: $\alpha_s C_F/\pi \;\; \rightarrow \;\; f_g(x,\mathbf{k}^2) = \mathcal{K}\, \partial_{(\ell n\, \mathbf{k}^2)}\, xg(x,\mathbf{k}^2)$
- Gap Survival Probability[2]: $S^2_{gap} \rightarrow$ 3% (5%) for LHC (Tevatron)
- Gluon radiation suppression[3]: Sudakov factor $S(\mathbf{k}^2, M_H^2) \sim \ell n^2\left(M_H^2/4\mathbf{k}^2\right)$
- Cutoff $\mathbf{k}_0^2$: Necessary to avoid infrared divergencies :: $\mathbf{k}_0^2 = 1\ \text{GeV}^2$.
- Electroweak vacuum expectation value: $v = 246$ GeV
- Gluon-proton form factor: $B = 5.5\ \text{GeV}^{-2}$

---

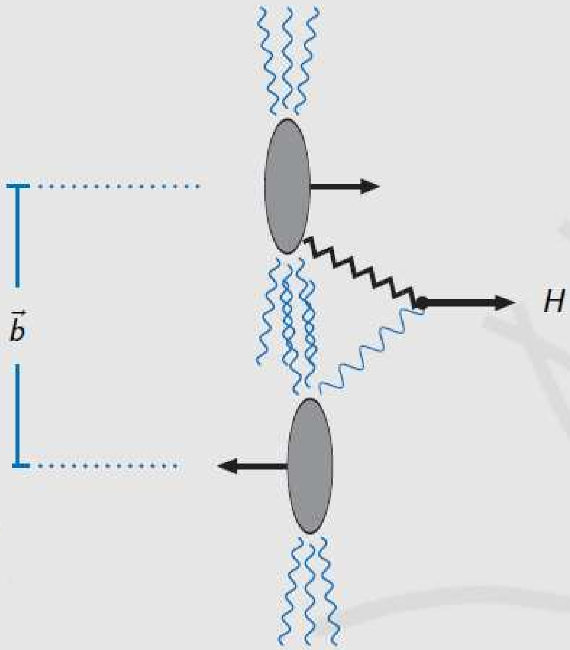
[1] Khoze, Martin, Ryskin, EJPC **14** (2000) 525

[2] Khoze, Martin, Ryskin, EJPC **18** (2000) 167

[3] Forshaw, `hep-ph/0508274`

Ultraperipheral Collisions

- Photon emission from the proton



$$\sigma(pp(A) \to p + H + p(A)) = 2\int_{\omega_0}^{\sqrt{s}/2} d\omega \, \frac{dn_i}{d\omega} \, \sigma_{\gamma p}(\omega, M_H),$$

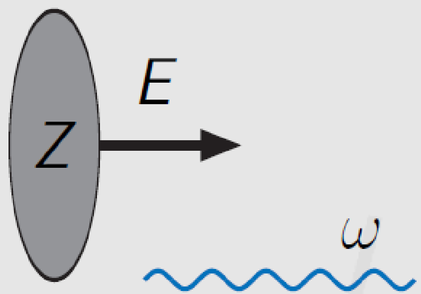
with photon fluxes

$$\frac{dn_p}{d\omega} = \frac{\alpha_{em}}{2\pi\omega}\left[1 + \left(1 - \frac{2\omega}{\sqrt{s}}\right)^2\right]\left(\ell nA - \frac{11}{6} + \frac{3}{A} - \frac{3}{2A^2} + \frac{1}{3A^2}\right).$$

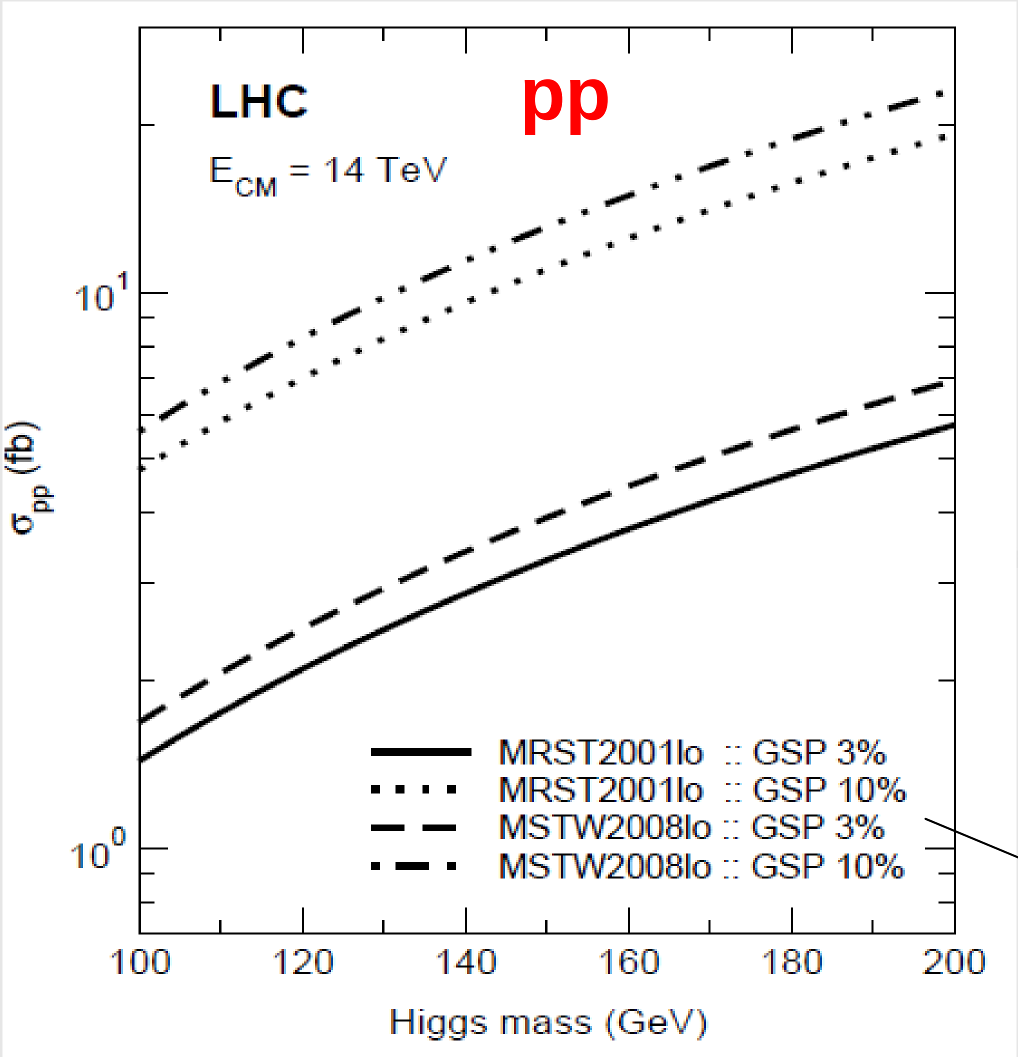
$$\frac{dn_A}{d\omega} = \frac{2Z^2\,\alpha_{em}}{\pi\omega}\left[\mu K_0(\mu)K_1(\mu) - \frac{\mu^2}{2}[K_1^2(\mu) - K_0^2(\mu)]\right].$$

- The photon virtuality obey the **Coherent condition** for its emission from a hadron under collision

$$Q^2 \lesssim 1/R^2$$



Photoproduction cross section



Subprocess	GSP (%)	σ_{pp} (fb)
$\mathbb{P}\mathbb{P}$	2.6	3.00
$\mathbb{P}\mathbb{P}$	0.4	0.47
$\gamma\gamma$	100	0.12
γp	3.0	1.77
γp	10.	5.92

M_H = 120 GeV

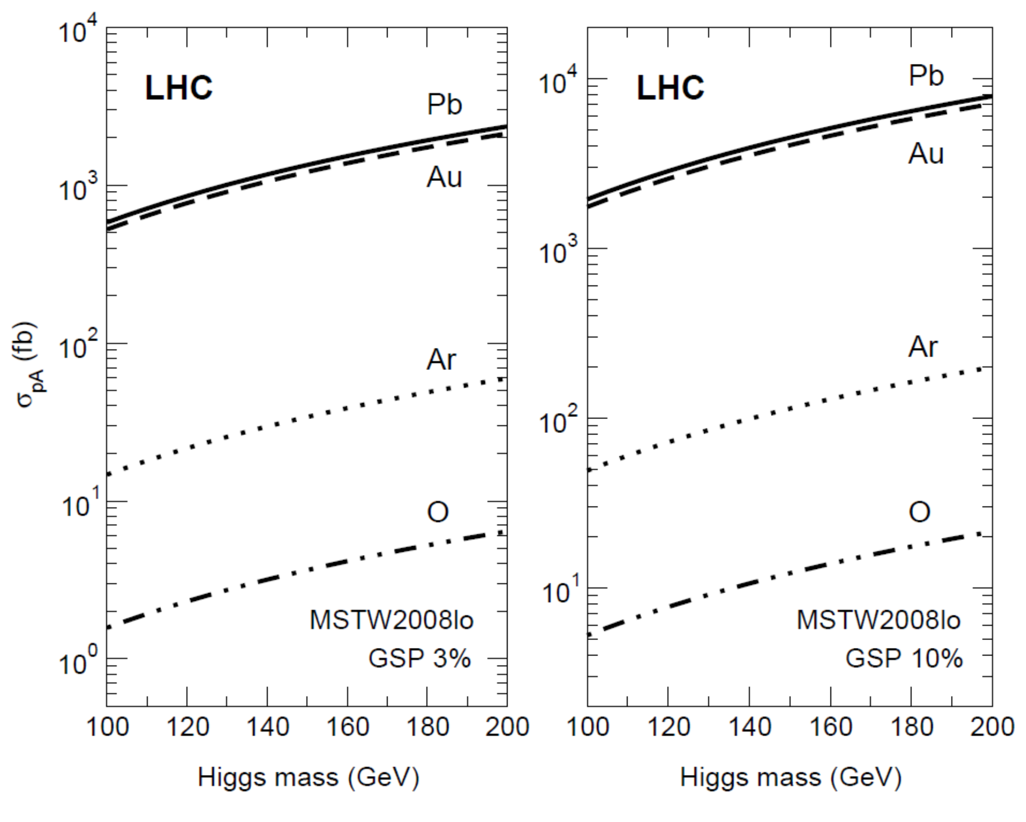
Cross section = 1.77-6 fb

Estimations for the GSP in the LHC energy

pA collisions

Process	σ (fb)	BR $\times\sigma$ (fb)	$\mathcal{L}$ (fb^{-1})	Events/yr
pp	1.77	1.27	1.(30.)	1 (30)
pp	5.92	4.26	1.(30.)	6 (180)
pPb	617.	444.	0.035	21
pPb	2056.	1480.	0.035	72

BR(H→bb-bar) = 72%



Conclusions

✓ GFPAE has been working in hard diffractive events

✓ Use of IS with absorptive corrections (gap survival probability)

→ describe Tevatron data for W^{+-} and Z^0 production

→ rate production for quarkonium + photon at LHC energies

$R^{(J/\psi)}_{SD}$ = 0,8 – 0,5 % $R^{(\Upsilon)}_{SD}$ = 0,6 – 0,4 %(first in literature)

→ predictions for heavy quark production (SD and DPE) at LHC energies possible to be verified in AA collision (diffractive cross section in pp, pA and AA collisions)

$$C\overline{C} \qquad B\overline{B}$$

A = Lead and Calcium

→ Higgs predictions in agreement with Hard Pomeron Exchange

Cross sections of Higgs production 1 fb (**DPE**); 60-80 fb (**SD**)

Conclusions

✓Exclusive photoproduction is promising for the LHC

→ strong suppression of backgrounds

→ cross section prediction ⟹ **2-6** fb

→ expecting between 1 and 6 events per year

→ additional signature with the Hγ associated production

→ High event rates for pA collisions

σ = 1 pb ⟹ pPb collisions

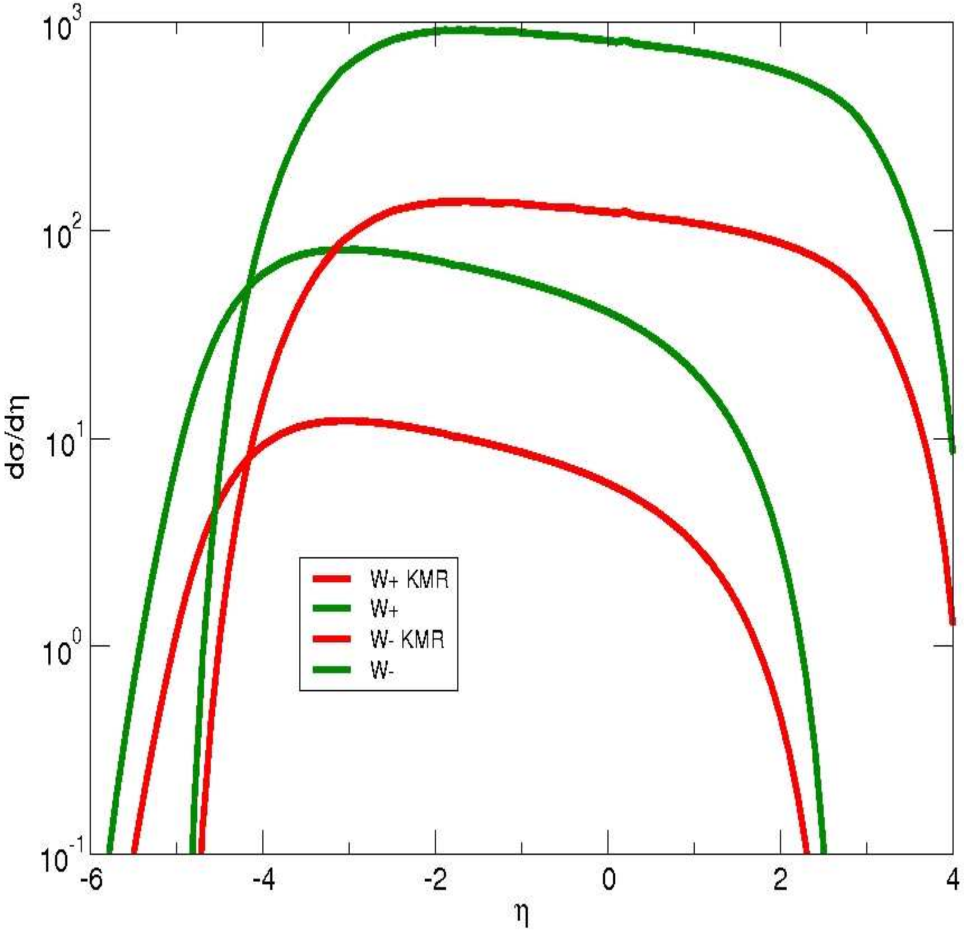
Next

DIFFRACTION IN NUCLEAR COLLISIONS

- ✓ Gap survival probability for nuclear collisions
- ✓ Dijets in hadronic and nuclear collisions
- ✓ ...

BACKUP

Predictions (LHC – 14 TeV)



High diffractive ratio

$$R_{KMR} = \frac{\int_{-1}^{1} \sigma_{diffractive}}{\int_{-1}^{1} \sigma_{inclusive}} = 0.311$$

Large range of pseudorapidity

$$-6 \leq \eta \leq 6$$

A. Bialas and W. Szeremeta, Phys. Lett. B 296, 191 (1992)

Bialas-Landshoff approach

Double Pomeron Exchange $p + p \rightarrow p + Q\overline{Q} + p$

$$\sigma_{I\!PI\!P}(\mathrm{BL}) = \frac{1}{2s\,(2\pi)^8} \int \overline{|M_{fi}|^2}\,[F(t_1)\,F(t_2)]^2\,dPH$$

$F(t)$ → nucleon form-factor

$$F(t) = \exp(b\,t) \qquad b = 2\ \mathrm{GeV}^{-2}$$

Differential phase-space factor

$$\begin{aligned} dPH &= d^4k_1\delta\left(k_1^2\right) d^4k_2\,\delta\left(k_2^2\right) d^4r_1\,\delta\left(r_1^2 - m_Q^2\right) \\ &\times\; d^4r_2\,\delta\left(r_2^2 - m_Q^2\right)\Theta\left(k_1^0\right)\Theta\left(k_2^0\right)\Theta\left(r_1^0\right)\Theta\left(r_2^0\right) \\ &\times\; \delta^{(4)}\left(p_1 + p_2 - k_1 - k_2 - r_1 - r_2\right), \end{aligned}$$

m_Q → mass of produced quarks

Bialas-Landshoff approach

Sudakov parametrization for momenta

$$Q = \frac{x}{s}p_1 + \frac{y}{s}p_2 + v, \qquad k_1 = x_1 p_1 + \frac{y_1}{s}p_2 + v_1,$$
$$k_2 = \frac{x_2}{s}p_1 + y_2 p_2 + v_2, \qquad r_2 = x_Q p_1 + y_Q p_2 + v_Q,$$

$v,\ v_1,\ v_2,\ v_Q$ → two-dimensional four-vectors describing the transverse component of the momenta

$p_1,\ p_2\ (k_1,\ k_2)$ → momenta for the incoming (outgoing) protons

$r_2\ (r_1)$ → momentum for the produced quark (antiquark)

Q → momentum for one of exchanged gluons

Bialas-Landshoff approach

Square of the invariant matrix element averaged over initial spins and summed over final spins

$$\overline{|M_{fi}|^2} = \frac{x_1 y_2 \, H}{(s x_Q y_Q)^2 \, (\delta_1 \delta_2)^{1+2\epsilon} \, \delta_1^{2\alpha' t_1} \delta_2^{2\alpha' t_2}} \left(1 - \frac{4\, m_Q^2}{s \delta_1 \delta_2}\right) \exp\left[2\beta\,(t_1 + t_2)\right]$$

$\delta_1 = 1 - x_1, \; \delta_2 = 1 - y_2, \; t_1 = -\vec{v}_1^2, \; t_2 = -\vec{v}_2^2$ $\qquad \beta = 1 \mathrm{GeV}^{-2}$

$\exp\left[2\beta\,(t_1 + t_2)\right]$ → effect of the momentum transfer dependence of the non-perturbative gluon propagator

$$H = S_{\mathrm{gap}}^2 \times 2s \left[\frac{4\pi m_Q \,(G^2 D_0)^3 \mu^4}{9\,(2\pi)^2}\right]^2 \left(\frac{\alpha_s}{\alpha_0}\right)^2$$

$\epsilon = 0.08, \; \alpha' = 0.25 \ \mathrm{GeV}^{-2}, \; \mu = 1.1 \ \mathrm{GeV}$

$G^2 D_0 = 30 \ \mathrm{GeV}^{-1} \mu^{-1}$

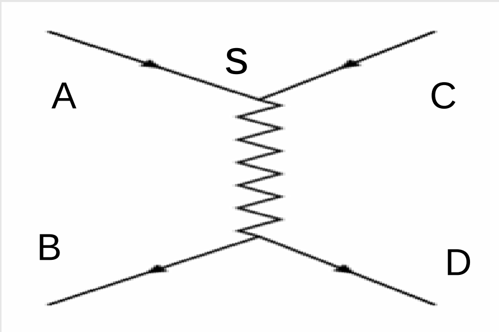
Processes in channels s and t

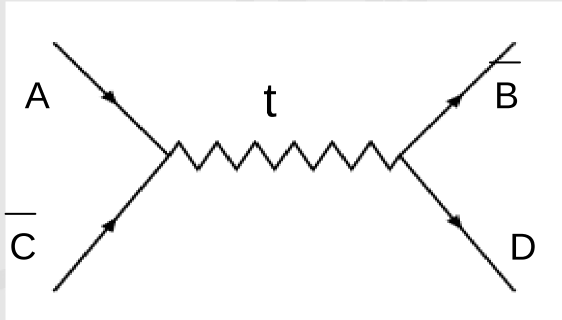
- Two body scattering can be calculated in terms of two independent invariants, s and t, Mandelstam variables

where $\begin{cases} s=(A+B)^2=(C+D)^2 \\ t=(A-C)^2=(B-D)^2 \end{cases}$

Square of center-of-mass energy

Square of the transfered four momentum





$$A_{AB\to CD}(s,t)=A_{A\bar{C}\to \bar{B}D}(t,s)$$ by crossing symmetry

$$A(s,t)\approx \frac{g^2}{m_\pi^2 - t}$$ pion exchange

g → coupling constant

Singularity (pole) in non-physical region t > 0 in s-channel diagram → $t = m_\pi^2$

Regge Theory

- At fixed t, with *s* >> *t*

- Amplitude for a process governed by the exchange of a trajectory α (t) is

$$A(s,t) \sim \left(s/s_0\right)^{\alpha(t)}$$

- No prediction for *t* dependence

- Elastic cross section

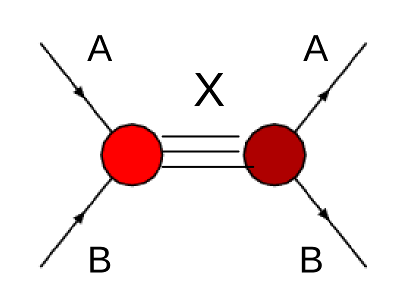
$$\frac{d\sigma_{el}}{dt} \sim s^{2\alpha(t)-2}$$

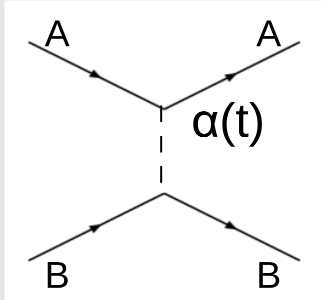
- Total cross section considering the optical theorem

Diffractive scattering

Consider elastic A B ⟶ A B

$$\frac{d\sigma_{el}}{dt} \approx \frac{1}{s^2}\sum_X \left| \text{[diagram]} \right|^2 \approx \frac{1}{s^2}\left| \text{[diagram]} \right|^2 \approx s^{2\alpha(t)-2}$$

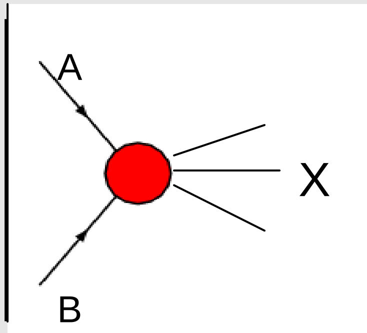


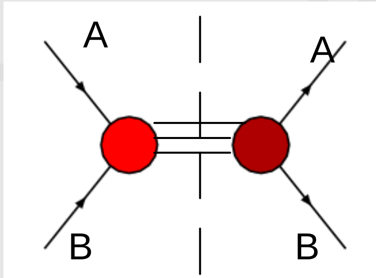


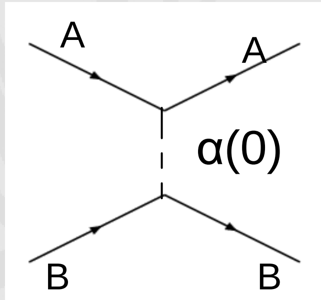
optical theorem

$$\sigma_{tot}^{AB} \approx \frac{1}{s}\,\mathrm{Im}\left(A_{el}^{AB}\right)_{t=o} \approx s^{\alpha(0)-1}$$

$$\sigma_{tot} = \frac{1}{2s}\sum_X \left| \text{[diagram]} \right| = \frac{1}{2s}\sum_X \text{[diagram]} \approx \frac{1}{s}\,\text{[diagram]}$$







by Regge

$\alpha(0) \approx 1+\varepsilon, \qquad \alpha(0) \leq 0.5$

Apparent contradiction

{ vacuum trajectory
Pomeron α_{IP} (t)
vacuum quantum numbers

Diffractive scattering

$$\alpha_{IP}(t) = 1.085 + 0.25t \qquad (p\,p,\ p\,\bar{p})$$

The interactions described by the exchange of a IP are called diffractive

so

$$\frac{d\sigma_{tot}^{AB}}{dt} \approx \frac{\beta_{AIP}^{2}(t)\,\beta_{BIP}^{2}(t)}{16\pi}\, s^{2\alpha_{IP}-2}$$

β_{iIP} ⟹ Pomeron coupling with external particles

Valid for $s \to \infty, \quad t/s \to 0$

High s ⟹ $\sigma_{tot}^{AB} \approx \beta_{AIP}(0)\,\beta_{BIP}(0)\, s^{\alpha_{IP}-1}$

Froissart limit

- No diffraction within a black disc

- It occurs only at periphery, b ~ R ⟹ in the Froissart regime, $R \propto \ln(s)$

- Unitarity demands

$$\sigma_{tot} \propto \sigma_{el} \propto \ln^2(s)$$
$$\sigma_{sd} \propto \ln(s) \ ,$$

⟹ i.e. $\sigma_{sd}/\sigma_{tot} \propto 1/\ln(s)$

- Donnachie-Landshoff approach ⟹ may not be distinguishable from logarithmic growth

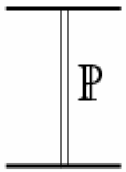
Any s^{λ} power behaviour would violate unitarity

↪ At some point should be modified by unitarity corrections

- Rate of growth ~ $s^{0.08}$ would violate unitarity only at large energies

Regge phenomenology in QCD

o Elastic amplitude ⇨ mediated by the Pomeron exchange



$$A_{el}(t) \propto \left[i - \text{ctg} \frac{\pi \alpha_{I\!P}(t)}{2} \right] \left(\frac{s}{s_0} \right)^{\alpha_{I\!P}(t)\, t)}$$

$$\alpha_{I\!P}(t) = \alpha^0_{I\!P} + \alpha'_{I\!P} t$$

What is the Pomeron?

o A Regge pole: not exactly, since $\alpha_{IP}(t)$ varies with Q^2 in DIS

o DGLAP Pomeron ⇨ specific ordering for radiated gluon

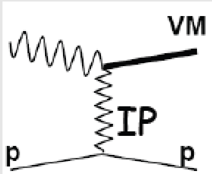
$$k^2_{i+1} < k^2_i \le Q^2 \quad \text{and} \quad x \le x_{i+1} \le x_i$$

o BFKL Pomeron ⇨ no ordering ⇨ no evolution in Q^2

o Other ideas?

Studies of diffraction

o In the beginning → hadron-hadron interactions

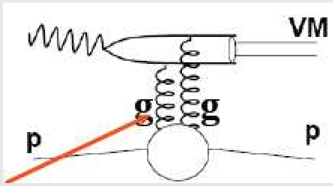


SOFT

low momentum transfer

o Exclusive diffractive production: ρ, φ, J/ψ, Y, γ

Gluon exchange



HARD

high momentum transfer

o Cross section

$$\sigma(W) \propto W^{\delta}$$

o δ expected to increase from soft (~ 0.2 is a "soft" Pomeron) to hard (~ 0.8 is a "hard" Pomeron)

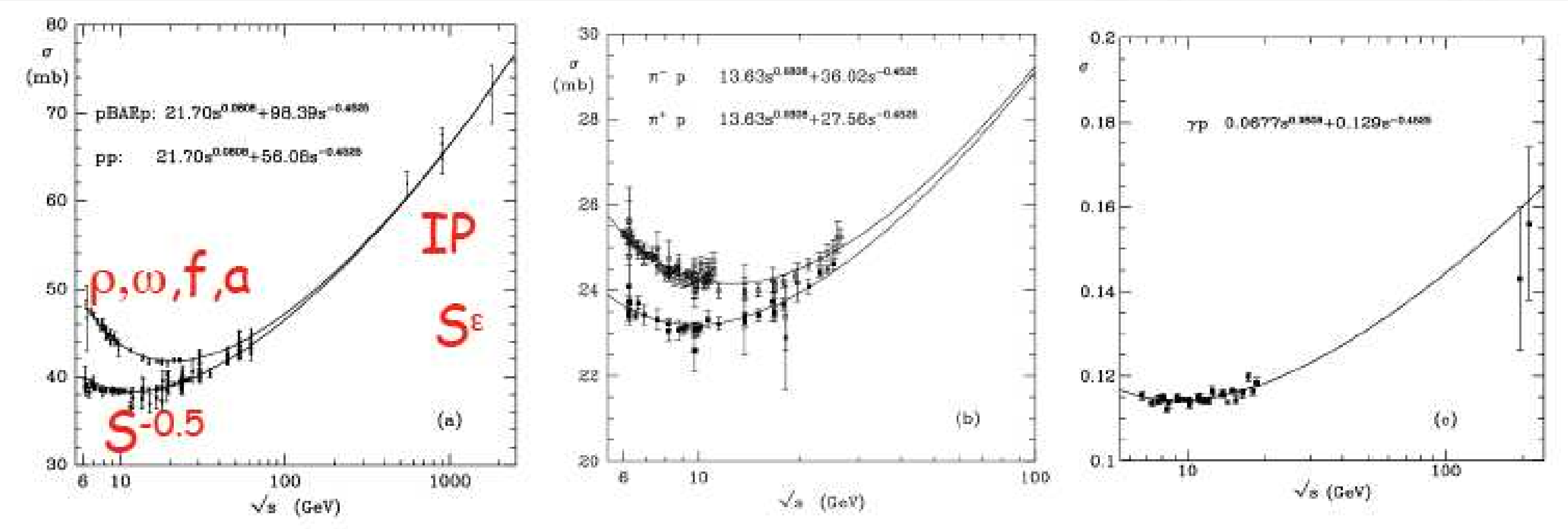
o Differential cross section

$$\frac{d\sigma}{dt} \propto e^{-b|t|}$$

Some results

✓ Many measurements in pp

✓ Pomeron exchange trajectory $\alpha(t) \sim 1.10 + 0.25\,t$



Pomeron universal and factorizable

applied to total, elastic, diffractive dissociation cross sections in *ep* collisions

Diffractive Structure Functions

✓ DDIS differential cross section can be written in terms of two structure functions

$$F_1^{D(4)} \quad \text{and} \quad F_2^{D(4)}$$

✓ Dependence of variables ⟶ $x, \; Q^2, \; x_{IP}, \; t$

✓ Introducing the longitudinal and transverse diffractive structure functions

$$F_L^{D(4)} = F_2^{D(4)} - 2xF_1^{D(4)} \qquad F_T^{D(4)} = 2xF_1^{D(4)}$$

✓ DDIS cross section is

$$\frac{d\sigma^D_{\gamma^* p}}{dx dQ^2 dx_{IP} dt} = \frac{4\pi\alpha_{em}^2}{xQ^4}\left\{1 - y + \frac{y^2}{2\left[1 + R^{D(4)}\left(x, Q^2, x_{IP}, t\right)\right]}\right\} F_2^{D(4)}\left(x, Q^2, x_{IP}, t\right)$$

✓ $R^{D(4)} = \dfrac{F_L^{D(4)}}{F_T^{D(4)}}$ is the longitudinal-to-transverse ratio

Diffractive Structure Functions

✓ Data are taken predominantly at small y

✓ Cross section ⟹ little sensitivy to $R^{D(4)}$

✓ $F_L^{D(4)} << F_T^{D(4)}$ for β < 0.8 – 0.9 ⟹ neglect $R^{D(4)}$ at this range

$$\frac{d\sigma^D_{\gamma^* p}}{dx\,dQ^2\,dx_{IP}\,dt} = \frac{4\pi\alpha_{em}^2}{xQ^4}\left(1 - y + \frac{y^2}{2}\right) F_2^{D(4)}\left(x, Q^2, x_{IP}, t\right)$$

✓ $F_2^{D(4)}$ ⟹ proportional to the cross section for diffractive γ*p scattering

$$F_2^{D(4)}\left(x, Q^2, x_{IP}, t\right) = \frac{Q^2}{4\pi\alpha_{em}^2} \frac{d\sigma^D_{\gamma^* p}}{dx_{IP}\,dt}$$

✓ $F_2^{D(4)}$ ⟹ dimensional quantity

$$F_2^{D(4)} \equiv \frac{dF_2^D(x, Q^2, x_{IP}, t)}{dx_{IP}\,dt}$$

F_2^D is dimensionless

Diffractive Structure Functions

✓ When the outgoing proton is not detected

no measurement of t

✓ Only the cross section integrated over t is obtained

$$\frac{d\sigma^{D}_{\gamma^* p}}{dx dQ^2 dx_{IP}} = \frac{4\pi\alpha_{em}^2}{xQ^4}\left(1 - y + \frac{y^2}{2}\right) F_2^{D(3)}\left(x, Q^2, x_{IP}\right)$$

✓ The structure function $F_2^{D(4)}$ is defined as

$$F_2^{D(3)}\left(x, Q^2, x_{IP}\right) = \int_0^\infty d\,|t|\, F_2^{D(4)}\left(x, Q^2, x_{IP}, t\right)$$

Diffractive Parton Distributions

✓ Factorization theorem holds for diffractive structure functions

✓ These can be written in terms of the diffractive partons distributions

✓ It represents the probability to find a parton in a hadron *h*, under the condition the *h* undergoes a diffractive scattering

✓ QCD factorization formula for F_2^D is

$$\frac{dF_2^D(x,Q^2,x_{IP},t)}{dx_{IP}dt} = \sum_i \int_x^{x_{IP}} d\xi \frac{df_i(\xi,\mu^2,x_{IP},t)}{dx_{IP}dt} \hat{F}_2^i\left(\frac{x}{\xi},Q^2,\mu^2\right)$$

✓ $df_i(\xi,\mu^2,x_{IP},t)/dx_{IP}dt$ is the diffractive distribution of parton *i*

✓ Probability to find in a proton a parton of type *i* carrying momentum fraction ξ

✓ Under the requirement that the proton remains intact except for a momentum transfer quantified by x_{IP} and *t*

Diffractive Parton Distributions

✓ Perturbatively calculable coefficients

$$\hat{F}_2^i\left(\frac{x}{\xi}, Q^2, \mu^2\right)$$

✓ Factorization scale ⟶ $\mu^2 = M^2$

✓ Diffractive parton distributions satisfy DGLAP equations

✓ Thus

$$\frac{\partial}{\partial \ln \mu^2} \frac{df_i(\xi, \mu^2, x_{IP}, t)}{dx_{IP} dt} = \sum_j \int_\xi^1 \frac{d\varsigma}{\varsigma} P_{ij}\left(\frac{\xi}{\varsigma}, \alpha_s(\mu)\right) \frac{df_j(\xi, \mu^2, x_{IP}, t)}{dx_{IP} dt}$$

✓ "**fracture function**" is a diffractive parton distribution integrated over t

$$\frac{df_i(\xi, \mu^2, x_{IP})}{dx_{IP}} = \int_{\frac{x_{IP}^2 m_N^2}{1 - x_{IP}}}^{\infty} d|t| \frac{df_i(\xi, \mu^2, x_{IP}, t)}{dx_{IP} dt}$$

Partonic Structure of the Pomeron

✓ It is quite usual to introduce a partonic structure for F_2^{IP}

✓ At Leading Order ➡ Pomeron Structure Function written as a superposition of quark and antiquark distributions in the Pomeron

$$F_2^{IP}(\beta, Q^2) = \sum_{q,\bar{q}} e_q^2 \beta q^{IP}(\beta, Q^2)$$

✓ $\beta = \frac{x}{x_{IP}}$ ➡ interpreted as the fraction of the Pomeron momentum carried by its partonic constituents

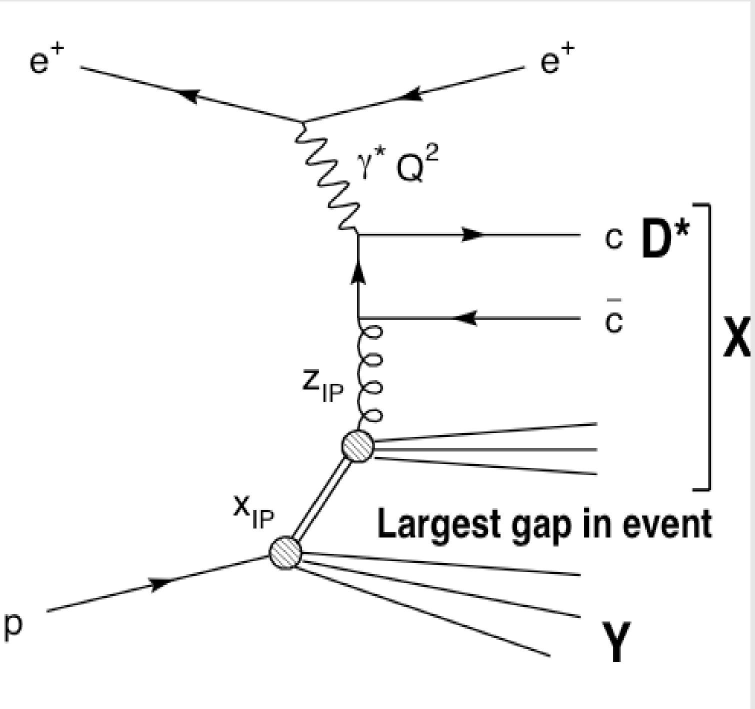
✓ $q^{IP}(\beta, Q^2)$ ➡ probability of find a quark q with momentum fraction β inside the Pomeron

✓ This interpretation makes sense only if we can specify unambigously the probability of finding a Pomeron in the proton and assume the Pomeron to be a real particle (Ingelman, Schlein, 1985)

Partonic Structure of the Pomeron

✓ Diffractive quark distributions and quark distributions of the Pomeron are related

$$\frac{df_q(\beta, Q^2, x_{IP}, t)}{dx_{IP}dt} = \frac{1}{16\pi^2} |g_{IP}(t)|^2 \, x_{IP}^{-2\alpha_{IP}(t)} q^{IP}(\beta, Q^2)$$



Representation of D* diffractive production in the infinite-momentum frame description of DDIS

- Introducing gluon distribution in the Pomeron

$$g^{IP}(\beta, Q^2)$$

- Related to $df_g / dx_{IP}dt$ by

$$\frac{df_g(\beta, Q^2, x_{IP}, t)}{dx_{IP}dt} = \frac{1}{16\pi^2} |g_{IP}(t)|^2 \, x_{IP}^{-2\alpha_{IP}(t)} g^{IP}(\beta, Q^2)$$

- At Next-to-Leading order, Pomeron Structure Function acquires a term containing $g^{IP}(\beta, Q^2)$

Diffractive processes

- Hadronic processes can be characterized by an energy scale

→ Soft processes – energy scale of the order of the hadron size (~ 1 fm)

pQCD is inadequate to describe these processes

$$\alpha_{soft}(t) = 1.08 + 0.25t$$

→ Hard processes – "hard" energy scale (> 1 GeV^2)

can use pQCD

"factorization theorems"

Separation of the perturbative part from non-perturbative

$$\alpha_{hard}(t) = 1.30 + 0.02t$$

- Most of diffractive processes at HERA → **"soft processes"**

Pomeron as composite

- Considering Regge factorization we have

$$F_2^{D(4)}\left(x,Q^2,x_{IP},t\right)= f_{IP/p}\left(x_{IP},t\right)F_2^{IP}(\beta,Q^2)$$

IP flux

IP Structure function

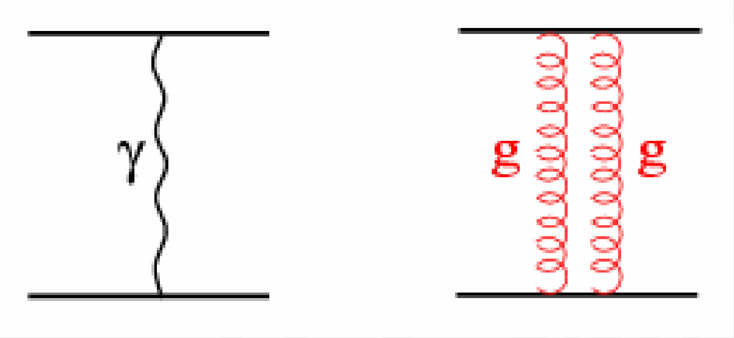
see MBGD & M. V. T. Machado 2001

Data → Good fit with added Reggeon for HERA

Pomeron as gluons

- Elastic amplitude ⇒ neutral exchange in t-channel
- Smallness of the real part of the diffractive amplitude ⇒ nonabeliance

Born graphs in the abelian and nonabelian (QCD) cases look like



The Pomeron

o From fitting elastic scattering data IP trajectory is much flatter than others

o For the intercept $\alpha'_{IP} \approx 0.25\ GeV^{-2}$ total cross sections implies $\alpha_{IP}(0) \approx 1$

o Pomeron ⟹ dominant trajectory in the elastic and diffractive processes

o Known to proceed via the exchange of **vacuum quantum numbers** in the *t*-channel

Regge-type $\alpha(t) = \alpha(0) + \alpha' t$

First measurements in h-h scattering

$$W^2 = (q+p)^2$$

$$\frac{d\sigma}{dt}(W) = \exp(b_0 t) W^{2[2\alpha_{IP}(t)+2]}$$

✓ $\alpha(0)$ and α' are fundamental parameters to represent the basic features of strong interactions

$$\frac{d\sigma}{dt}(W) = W^{4\alpha(0)-4}\exp(bt)$$

$$b = b_0 + 4\alpha' \ln(W)$$

✓ α' ⟹ energy dependence of the transverse system

Pomeron structure function

- Pomeron structure function has been modeled in terms of a light flavor singlet distribution Σ(z)

- Consists of u, d and s quarks and antiquarks and a gluon distribution g(z)

- z is the longitudinal momentum fraction of the parton entering the hard subprocess with respect of the diffractive exchange

- (z = β) for the lowest order quark-parton model process and 0 < β < z for higher order processes

- Quark singlet and gluon distributions are parametrized at Q^2_0

$$zf_{i/IP}(z, Q_0^2) = A_i z^{B_i} (1-z)^{C_i} \exp\left[-\frac{0.01}{(1-z)}\right]$$

Pomeron structure function

- Experimental determination of the diffractive PDFs involves the following cuts

$$\beta < 0.8, M_X > 2GeV; Q^2 < 8.5GeV^2$$

- Quark singlet distribution, data requires inclusion of parameters A_q, B_q and C_q

- Gluon density is weakly constrained by data which are found to be insensitive to the B_g parameter

- **FIT A** - Gluon density is parametrized using only A_g and C_g parameters (Q^2_0 = 1.75 GeV^2)

- This procedure is not sensitive to the gluon PDF and a new adjustment was done with C_g = 0

- **FIT B** - Gluon density is a simple constant at the starting scale for evolution (Q^2_0 = 2.5 GeV^2)

Pomeron structure function

Parameter	*Value*
α'_{IP}	$0.06^{+0.19}_{-0.06} GeV^{-2}$
B_{IP}	$5.5^{+2.0}_{-0.7} GeV^{-2}$
$\alpha_{IR}(0)$	0.50 ± 0.10
α'_{IR}	$0.3^{+0.6}_{-0.3} GeV^{-2}$
B_{IR}	$1.6^{+1.6}_{-0.4} GeV^{-2}$
m_c	$1.4 \pm 0.2 GeV$
m_b	$4{,}5 \pm 0.5 GeV$
$\alpha_8^{(5)}(M_Z^2)$	0.118 ± 0.002

- Values of fixed parameters (masses) and their uncertainties, as used in the QCD fits.
- α'_{IP} and B_{IP} (strongly anti-correlated) are varied simultaneously to obtain the theoretical errors on the fits (as well as α'_{IR} and B_{IR}).
- Remaining parameters are varied independently.
- Theoretical uncertainties on the free parameters of the fit are sensitive to the variation of the parametrization scale Q^2_0